



REFERENCES

1. Tabor, M., Chaos and Integrability in Nonlinear Dynamics, John-Wiley & Sons, New York, 1989.
2. Bertrand, J., J. Math., (i), XVII, 121, 1852.
3. Dorizzi, B., B. Grammaticos, and A. Ramani, "A New Class of Integrable Systems," J. Math. Phys., 24, 2282-2288, 1983.
4. Grammaticos, B., B. Dorizzi, and A. Ramani, "Integrability of Hamiltonians with Third and Fourth Degree Polynomial Potentials," J. Math. Phys., 24, 2289-2295, 1983.
5. Dorizzi, B., B. Grammaticos, A. Ramani, and P. Winternitz, "Integrable Hamiltonian Systems with Velocity dependent potentials," J. Math. Phys., 26, 3070-3079, 1985.
6. Holt, C. R., "Construction of new integrable Hamiltonians in Two Degrees of Freedom," J. Math. Phys., 23, 1037-1046, 1982.
7. Ablowitz, M. J., A. Ramani, and H. Segur, "A Connection Between Nonlinear Evolution Equations and Ordinary Differential Equations," J. Math. Phys., 21, 715-721, 1980.

8. McLeod, J. B., and P. J. Olver, "The Connection Between Partial Differential Equations Soluble by Inverse Scattering and Ordinary Differential Equations of Painleve Type," Siam J. Math. Anal., 14, 488-456, 1983.
9. Ablowitz, M. J., and H. Segur, "Exact Linearization of a Painlevé Transcendent," Phys. Rev. Lett., 38, 1103-1106, 1977.
10. Toda, M., J. Phys. Soc. Jpn., 23, 501, 1967; Progr. Theor. Phys. Suppl., 45, 174, 1970.
11. Flaschka, H., "On the Toda Lattice II," Prog. Theor. Phys., 51, 703-716, 1974.
12. Ramani, A., B. Grammaticos, and T. Bountis, "The Painlevé Property and Singularity Analysis of Integrable and Non-Integrable Systems," Phys. Rep., 180, 159-245, 1989.
13. Grammaticos, B., B. Dorizzi, and R. Pedjen, "Painlevé Property and Integrals of Motion for the Hénon-Heiles System," Phys. Lett., 89A, 111-113, 1982.
14. Bountis, T., H. Segur, and F. Vivaldi, "Integrable Hamiltonian Systems and the Painlevé Property," Phys. Rev. Lett., 25, 1257-1264, 1982.
15. Green, J., quoted in Ref.[7].

16. Chang, Y. F., M. Tabor, and J. Weiss, "Analytic structure of the Hénon-Heiles Hamiltonian in Integrable and Nonintegrable Regimes," J. Math. Phys., 23, 531-538, 1982.
17. Dorizzi, B., B. Grammaticos, R. Padjen, and V. Papageorgiou, "Integrals of Motion for Toda Systems with Unequal Masses," J. Math. Phys., 25, 2200-2211, 1984.
18. Ramani, A., quoted in Ref.[7].
19. Ramani, A., B. Dorizzi, and B. Grammaticos, "Painlevé Conjecture Revisited" Phys. Rev. Lett., 49, 1539-1541, 1982.
20. Ramada, A. F., A. Ramani, B. Dorizzi, and B. Grammaticos, "The Weak Painlevé Property as a Criterion for the Integrability of Dynamical Systems," J. Math. Phys., 26, 708-710, 1985.
21. Grammaticos, B., B. Dorizzi, and A. Ramani, "Hamiltonians with Higher-Order Integrals and the Weak Painlevé Concept," J. Math. Phys., 25, 3470-3473, 1984.
22. Weiss, J., M. Tabor, and J. Carnevale, J. Math. Phys., 24, 522, 1983.



APPENDIX.

DIRECT SEARCH OF THE FREE-END TODA LATTICE

In this appendix we use a direct calculation method at order 6 in the velocities for the cases of the free-end Toda lattice.

$$H = \frac{p_1^2}{2m_1} + \frac{p_2^2}{2m_2} + \frac{p_3^2}{2m_3} + e^{\epsilon(q_1 - q_2)} + e^{q_2 - q_3}. \quad (A1)$$

We introduce the following change of variables:

$$x = \epsilon(q_1 - q_2), \quad y = q_2 - q_3, \quad z = m_1 q_1 + m_2 q_2 + m_3 q_3. \quad (A2)$$

The equations of motion are derived from the Lagrange equations:

$$\begin{aligned} \frac{\partial}{\partial t}(m_1 \dot{q}_1) &= -\epsilon X, & \frac{\partial}{\partial t}(m_2 \dot{q}_2) &= \epsilon X - Y, \\ \frac{\partial}{\partial t}(m_3 \dot{q}_3) &= Y, \end{aligned} \quad (A3)$$

$$\text{with } X = e^{\epsilon(q_1 - q_2)} = e^x, \quad Y = e^{q_2 - q_3} = e^y.$$

From Eq. (A2), it is obvious that

$$\begin{aligned} \ddot{x} &= \epsilon(\ddot{q}_1 - \ddot{q}_2) = \frac{\epsilon}{m_2} [Y - \epsilon(\frac{m_1 + m_2}{m_1})X], \\ \ddot{y} &= \ddot{q}_2 - \ddot{q}_3 = \frac{\epsilon}{m_2} [X - (\frac{m_2 - m_3}{\epsilon m_3})Y], \\ \ddot{z} &= m_1 \ddot{q}_1 + m_2 \ddot{q}_2 + m_3 \ddot{q}_3 = 0, \end{aligned} \quad (A4)$$

which, after a scaling in time, read

$$\ddot{x} = Y - \alpha X, \quad \ddot{y} = X - \beta Y, \quad (A5)$$

$$\text{with } \alpha = \frac{\epsilon(m_1 + m_2)}{m_1}, \quad \beta = \frac{(m_2 + m_3)}{\epsilon m_3}.$$

Let us consider the constant of motion of order 6 in the velocities.

$$\begin{aligned} C = & e_0 \dot{x}^6 + e_1 \dot{x}^5 \dot{y} + e_2 \dot{x}^4 \dot{y}^2 + e_3 \dot{x}^3 \dot{y}^3 + e_4 \dot{x}^2 \dot{y}^4 + e_5 \dot{x} \dot{y}^5 \\ & + e_6 \dot{y}^6 + f_0 \dot{x}^4 + f_1 \dot{x}^3 \dot{y} + f_2 \dot{x}^2 \dot{y}^2 + f_3 \dot{x} \dot{y}^3 + f_4 \dot{y}^4 + g_0 \dot{x}^2 \\ & + g_1 \dot{x} \dot{y} + g_2 \dot{y}^2 + h. \end{aligned}$$

The condition of the constancy of C can be written as

$$\begin{aligned} 0 = dC = & 6e_0 \ddot{x}^5 + 5e_1 \ddot{x}^4 \dot{y} + e_1 \ddot{x}^5 \dot{y} + 4e_2 \ddot{x}^3 \dot{x} \dot{y}^2 + 2e_2 \ddot{x}^4 \dot{y} \dot{y} \\ & + 3e_3 \ddot{x}^2 \dot{x} \dot{y}^3 + 3e_3 \ddot{x}^3 \dot{y}^2 \dot{y} + 2e_4 \ddot{x} \dot{x} \dot{y}^4 + 4e_4 \ddot{x}^2 \dot{y}^3 + e_5 \ddot{x} \dot{y}^5 + 5e_5 \ddot{x} \dot{y} \dot{y} \\ & + 6e_6 \ddot{y}^5 + \frac{\partial f_0}{\partial x} \dot{x}^5 + \frac{\partial f_0}{\partial y} \dot{x}^4 \dot{y} + 4f_0 \dot{x}^3 \dot{x} \dot{y} + \frac{\partial f_1}{\partial x} \dot{x}^4 \dot{y} + \frac{\partial f_1}{\partial y} \dot{x}^3 \dot{y}^2 \\ & + 3f_1 \dot{x}^2 \dot{y} \dot{x} + f_1 \dot{x}^3 \dot{y} \dot{y} + \frac{\partial f_2}{\partial x} \dot{x}^3 \dot{y}^2 + \frac{\partial f_2}{\partial y} \dot{x}^2 \dot{y}^3 + 2f_2 \dot{x} \dot{x} \dot{y} \\ & + 2f_2 \dot{x}^2 \dot{y} \dot{y} + f_3 \dot{x} \dot{y}^3 + 3f_3 \dot{x} \dot{y}^2 \dot{y} + \frac{\partial f_3}{\partial x} \dot{x}^2 \dot{y}^3 + \frac{\partial f_3}{\partial y} \dot{x} \dot{y}^4 + \frac{\partial f_4}{\partial x} \dot{x} \dot{y}^4 \\ & + \frac{\partial f_4}{\partial y} \dot{y}^5 + 4f_4 \dot{y}^3 \dot{y} + \frac{\partial g_0}{\partial x} \dot{x}^3 + \frac{\partial g_0}{\partial y} \dot{x}^2 \dot{y} + 2g_0 \dot{x} \dot{x} + g_1 \dot{x} \dot{y} \\ & + g_1 \dot{x} \dot{y} + \frac{\partial g_1}{\partial x} \dot{x}^2 \dot{y} + \frac{\partial g_1}{\partial y} \dot{x} \dot{y}^2 + \frac{\partial g_2}{\partial x} \dot{x} \dot{y}^2 + \frac{\partial g_2}{\partial y} \dot{y}^3 + 2g_2 \dot{y} \dot{y} + \frac{\partial h}{\partial x} \\ & + \frac{\partial h}{\partial y}. \end{aligned} \quad (A6)$$

We can restrict ourselves to constant e_i 's

Regrouping and equating to zero the coefficients of each monomial in the velocities, we obtain at order 5.

$$\begin{aligned} 6e_0 \ddot{x} + e_1 \ddot{y} + \frac{\partial f_0}{\partial x} &= 0, & 5e_1 \ddot{x} + 2e_2 \ddot{y} + \frac{\partial f_0}{\partial y} + \frac{\partial f_1}{\partial x} &= 0, \\ 4e_2 \ddot{x} + 3e_3 \ddot{y} + \frac{\partial f_1}{\partial y} + \frac{\partial f_2}{\partial x} &= 0, & 3e_3 \ddot{x} + 4e_4 \ddot{y} + \frac{\partial f_2}{\partial y} + \frac{\partial f_3}{\partial x} &= 0, \\ 2e_4 \ddot{x} + 5e_5 \ddot{y} + \frac{\partial f_3}{\partial y} + \frac{\partial f_4}{\partial x} &= 0, & e_5 \ddot{x} + 6e_6 \ddot{y} + \frac{\partial f_4}{\partial y} &= 0. \end{aligned} \quad (A7)$$

The integrability condition for f_L reads

$$\begin{aligned} \frac{\partial^5}{\partial y^5}(6e_0 \ddot{x} + e_1 \ddot{y}) - \frac{\partial^5}{\partial x^4 \partial y}(5e_1 \ddot{x} + 2e_2 \ddot{y}) + \frac{\partial^5}{\partial x^2 \partial y^3}(4e_2 \ddot{x} + 3e_3 \ddot{y}) \\ - \frac{\partial^5}{\partial x^3 \partial y^2}(3e_3 \ddot{x} + 4e_4 \ddot{y}) + \frac{\partial^5}{\partial x^4 \partial y}(2e_4 \ddot{x} + 5e_5 \ddot{y}) - \frac{\partial^5}{\partial x^5}(e_5 \ddot{x} + 6e_6 \ddot{y}) = 0. \end{aligned} \quad (A8)$$

In the particular case where x and y are given by Eq.(A5),
Eq.(A8) reduces to

$$\frac{\partial^5}{\partial y^5}(6e_0 \ddot{x} + e_1 \ddot{y}) - \frac{\partial^5}{\partial x^5}(e_5 \ddot{x} + 6e_6 \ddot{y}) = 0;$$

Thus $6e_0 = e_1 \beta$, (A9.1)

$$\alpha e_5 = 6e_6. \quad (A9.2)$$

The integration for the f_i is

$$f_0 = e_1(\alpha\beta - 1)X + (2e_2\beta - 5e_1)Y,$$

$$\equiv A_0 X + B_0 Y,$$

$$f_1 = (5e_1\alpha - 2e_2)X + (3e_3\beta - 4e_2)Y,$$

$$\equiv A_1 X + B_1 Y,$$

$$f_2 = (4e_2\alpha - 3e_3)X + (4e_4\beta - 3e_3)Y,$$

$$\equiv A_2 X + B_2 Y,$$

$$f_3 = (3e_3\alpha - 4e_4)X - 2e_4 Y,$$

$$\equiv A_3 X + B_3 Y,$$

$$f_4 = 2e_4\alpha X + (6e_6\beta - e_5)Y,$$

$$\equiv A_4 X + B_4 Y. \quad (A10)$$

At third order we obtain

$$\begin{aligned} 4f_0 \ddot{x} + f_1 \ddot{y} + \frac{\partial g_0}{\partial x} &= 0, \\ 3f_1 \ddot{x} + 2f_2 \ddot{y} + \frac{\partial g_0}{\partial y} + \frac{\partial g_1}{\partial x} &= 0, \\ 2f_2 \ddot{x} + 3f_3 \ddot{y} + \frac{\partial g_1}{\partial y} + \frac{\partial g_2}{\partial x} &= 0, \\ f_3 \ddot{x} + 4f_4 \ddot{y} + \frac{\partial g_2}{\partial y} &= 0. \end{aligned} \quad (A11)$$

The integrability condition for g_1 reads

$$\begin{aligned} & -\frac{\partial^3}{\partial x^3}(f_3 \ddot{x} + 4f_4 \ddot{y}) + \frac{\partial^3}{\partial x^2 \partial y}(2f_2 \ddot{x} + 3f_3 \ddot{y}) - \frac{\partial^3}{\partial x \partial y^2}(3f_1 \ddot{x} + 2f_2 \ddot{y}) \\ & + \frac{\partial^3}{\partial y^3}(4f_0 \ddot{x} + f_1 \ddot{y}) = 0. \end{aligned} \quad (\text{A12})$$

Eq. (A12) is an identity in terms of the independent functions x^2, XY, Y^2 , in this case. We obtain a system in terms of A_i, B_i, α, β :

$$4B_0 - \beta B_1 = 0, \quad (\text{A13.1})$$

$$\alpha A_3 - 4A_4 = 0, \quad (\text{A13.2})$$

$$\begin{aligned} & (\beta A_1 - B_1) + 2(\alpha B_2 - A_2) + 3(\beta A_3 - B_3) + 4(\alpha B_0 - A_0) \\ & = (\alpha B_3 - A_3) + 2(\beta A_2 - B_2) + 3(\alpha B_1 - A_1) + 4(\beta A_4 - B_4). \end{aligned} \quad (\text{A13.3})$$

It is possible to calculate g :

$$\begin{aligned} g_0 &= \frac{1}{2}(4A_0\alpha - A_1)x^2 + [4(B_0\alpha - A_0) + (A_1\beta - B_1)]XY \\ &\quad + \frac{1}{2}[2\beta B_2 - 3B_1]Y^2, \\ &\equiv C_0 X^2 + D_0 XY + E_0 Y^2, \end{aligned}$$

$$\begin{aligned} g_1 &= \frac{1}{2}(3\alpha A_1 - 2A_2)x^2 + [3(\alpha B_1 - A_1) + 2(\beta A_2 - B_2) \\ &\quad + 4(A_0 - \alpha B_0) - (A_1\beta - B_1)]XY + (3\beta B_3 - 2B_2)Y^2, \\ &\equiv C_1 X^2 + D_1 XY + E_1 Y^2, \end{aligned}$$

$$\begin{aligned} g_2 &= (\alpha A_2 - \frac{3}{2}A_3)x^2 + (4\beta A_4 + \alpha B_3 - A_3 - 4B_4)XY \\ &\quad + (2\beta B_4 - 4B_4)Y^2, \\ &\equiv C_2 X^2 + D_2 XY + E_2 Y^2. \end{aligned} \quad (\text{A14})$$

At first order we obtain

$$\begin{aligned} 2g_0 \ddot{x} + g_1 \ddot{y} + \frac{\partial h}{\partial x} &= 0, \\ g_1 \ddot{x} + 2g_2 \ddot{y} + \frac{\partial h}{\partial y} &= 0, \end{aligned} \quad (A15)$$

The compatibility condition for the last equation reads

$$\frac{\partial}{\partial y} (2g_0 \ddot{x} + g_1 \ddot{y}) = \frac{\partial}{\partial x} (g_1 \ddot{x} + 2g_2 \ddot{y}). \quad (A16)$$

Eq.(A16) is an identity in terms of the independent functions $x^3, y^3, x^2 y, xy^2$, in this case .We obtain a system in terms of C_i, D_i, α, β :

$$4(D_0 - \alpha E_0) - 2(\beta D_1 - E_1) - (D_1 - \alpha E_1) - 2(E_2 - \beta D_2) = 0, \quad (A17.1)$$

$$2(C_0 - \alpha D_0) + (D_1 - \beta C_1) + 2(\alpha D_1 - C_1) + 4(\beta C_2 - D_2) = 0, \quad (A17.2)$$

$$2E_0 - \beta E_1 = 0, \quad (A17.3)$$

$$2C_2 - \alpha C_1 = 0. \quad (A17.4)$$

Eqs. (A13.1), (A13.2), (A17.3) and (A17.4) reduced to the equation in terms α and β :

$$(3\alpha\beta - 4)(-3\alpha\beta + 6)(\alpha\beta - 4) = 0.$$

a) $\alpha = 2, \beta = 2$: This integral is the product of the constant of degree 3 in the velocities.

b) $\alpha = 2, \beta = 1$: This integral is the product of the constant of degree 4 in the velocities.

c) $\alpha = \frac{2}{3}, \beta = 2$: This corresponds to

$$m_1 = 3\epsilon \frac{(2\epsilon - 1)}{2 - 3\epsilon}, \quad m_2 = 2\epsilon - 1, \quad m_3 = 1, \quad 1 < \epsilon < 2.$$

Substitute Eqs. (A13.1), (A13.2), (A17.3) and $\alpha = \frac{2}{3}, \beta = 2$ into Eq. (A13.3), which gives

$$-e_5 + 12e_6 = 0, \quad (\text{A18})$$

But Eq. (A9) is $4e_5 - 36e_6 = 0.$

This means that $e_5 = 0, e_6 = 0.$

For case $\alpha = \frac{2}{3}, \beta = 2$, we find the functions f_i 's, g_i 's and h :

$$\begin{aligned} f_0 &= 4X - 8Y, & f_1 &= 14X - 16Y, & f_2 &= \frac{50}{3}X - 10Y, \\ f_3 &= 8X - 2Y, & f_4 &= \frac{4}{3}X, \\ g_0 &= -\frac{5}{3}X^2 + \frac{20}{3}XY + 4Y^2, \\ g_1 &= -\frac{8}{3}X^2 + 6XY + 4Y^2, \\ g_2 &= -\frac{8}{9}X^2 + \frac{4}{3}XY + Y^2. \\ h &= \frac{4}{9}X^2Y + \frac{4}{27}X^3. & (e_4 &= 1) \end{aligned}$$

The constant C is given by

$$\begin{aligned} C &= 4\dot{x}^6 + 12\dot{x}^5\dot{y} + 13\dot{x}^4\dot{y}^2 + 6\dot{x}^3\dot{y}^3 + \dot{x}^2\dot{y}^4 + 4(e^x - 2e^y)\dot{x}^4 \\ &\quad + (14e^x - 16e^y)\dot{x}^3\dot{y} + 10(\frac{5}{3}e^x - e^y)\dot{x}^2\dot{y}^2 + 2(4e^x - e^y)\dot{x}\dot{y}^3 \\ &\quad + \frac{4}{3}e^x\dot{y}^4 + (-\frac{5}{3}e^{2x} + \frac{20}{3}e^{x+y} + 4e^{2y})\dot{y}^2 + \frac{4}{27}e^{3x} + \frac{4}{9}e^{2x+y} \\ &\quad + (-\frac{8}{3}e^{2x} + 6e^{x+y} + 4e^{2y})\dot{x}\dot{y} + (-\frac{8}{9}e^{2x} + \frac{4}{3}e^{x+y} + e^{2y})\dot{y}^2. \end{aligned}$$

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