

CHAPTER III

PRINCIPLES FOR ANALYZING GENERAL FLUID NETWORKS

3.1 Introduction

A general fluid network is formulated from a number of n nodes, each of which is identified by an index such as i or j .

The type of node i is specified as T_i , which assumes one of the following integral values:

- 0 - Pressure unspecified at node i .
- 1 - Pressure specified at node i .
- 2 - Injection or withdrawal rate specified at node i .
- 3 - Terminal node i with a specified injection or withdrawal rate.

The nodal connections are connected directly to other nodes by a pipeline or equipment. A connection matrix is established with possible values for a representative element C_{ij} , as follows:

- 1 - Node i and node j are joined by a pipeline.
- 2 - Centrifugal pump that pumps from node i to node j .
- 3 - Centrifugal compressor that compresses from node i to node j .

For nodal connection $C_{ij} = 1$, the pipeline diameter D_{ij} , length L_{ij} , and roughness ε_{ij} , are symmetrical joining node i and node j .

Nodes at which the specified injection rate is represented positive value or withdrawal rate as negative value.

Note:

For nodal connection across an equipment such as pump or compressor, if $C_y = 2$ or $C_y = 3$ then always $C_{ji} = 0$ because it can not operate in reverse flow.

The program always considers node i as the receiving node. Therefore, the flow rate within pipeline connection given as " Q_{ji} " is positive value for flow from node j to node i and negative for the reverse direction.

3.2 Flow in Pipelines

3.2.1 For Liquid

The flow rate from node j to node i is given by:

$$Q_{ji} = \sqrt{y / (\alpha_{ji} f_{F_{ji}})} \quad \text{for } y > 0 \quad (3.1)$$

The flow rate from node i to node j is given by:

$$Q_{ij} = -\sqrt{-y / (\alpha_{ji} f_{F_{ji}})} \quad \text{for } y < 0 \quad (3.2)$$

Here:

$$\alpha_{ji} = \frac{32\rho L_{ji}}{\pi^2 D_{ji}^5}, \quad \beta = \rho g \quad (3.3)$$

$$y = p_j - p_i + \beta(z_j - z_i) \quad (3.4)$$

3.2.2 For Gas

Inclined Flow ($z_i \neq z_j$):

The flow rate from node j to node i is given by:

$$Q_{sc-ji} = \frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{w}{\delta_{ji}(\phi_{ji} - 1)}} \quad \text{for } w > 0 \quad (3.5)$$

The flow rate from node i to node j is given by:

$$Q_{sc-ij} = -\frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{-w}{\delta_{ji}(\phi_{ji} - 1)}} \quad \text{for } w < 0 \quad (3.6)$$

Here:

$$w = p_j^2 - \phi_{ji} p_i^2 \quad (3.7)$$

$$\lambda_{ji} = \left(\frac{M}{Z_{avg} RT} \right) \left(\frac{\pi D_{ji}^2}{4} \right) \quad (3.8)$$

$$\delta_{ji} = \left(\frac{2f_{Fj} L_{ji}}{g D_{ji} (z_i - z_j)} \right) \quad (3.9)$$

$$\phi_{ji} = \exp \left(\frac{2Mg(z_i - z_j)}{Z_{avg} RT} \right) \quad (3.10)$$

$$\psi_{sc} = \frac{p_{sc} M}{Z_{sc} RT_{sc}} \quad (3.11)$$

Horizontal Flow ($z_i = z_j$):

The flow rate from node j to node i is given by:

$$Q_{sc-ji} = \frac{A_{ji}}{\psi_{sc}} \sqrt{\xi_{ji} (p_j^2 - p_i^2)} \quad \text{for } p_j > p_i \quad (3.12)$$

The flow rate from node i to node j is given by:

$$Q_{sc-ij} = -\frac{A_{ji}}{\psi_{sc}} \sqrt{-\xi_{ji} (p_j^2 - p_i^2)} \quad \text{for } p_j < p_i \quad (3.13)$$

Here:

$$A_{ji} = \frac{\pi D_{ji}^2}{4} \quad (3.14)$$

$$\xi_{ji} = \left(\frac{M}{Z_{avg} RT} \right) \left(\frac{D_{ji}}{4f_{Fji} L_{ji}} \right) \quad (3.15)$$

$$\psi_{sc} = \frac{p_{sc} M}{Z_{sc} RT_{sc}} \quad (3.16)$$

3.3 Flow in Equipment

3.3.1 For Liquid

There are two separate cases to be considered for each of three possibilities as follows:

1. $Q_{ji} > 0$ for flow across the pump from node j to node i:

$$Q_{ji} = 0 \quad \text{for } p_i + \beta z_i > p_j + \beta z_j + a_{ji} \quad (3.17)$$

$$Q_{ji} = \sqrt{a_{ji}/b_{ji}} \quad \text{for } p_j + \beta z_j > p_i + \beta z_i \quad (3.18)$$

$$Q_{ji} = \sqrt{(p_j - p_i + a_{ji} + \beta(z_j - z_i))/b_{ji}} \quad \text{otherwise} \quad (3.19)$$

2. $Q_{ij} < 0$ for flow across the pump from node i to node j:

$$Q_{ij} = 0 \quad \text{for } p_j + \beta z_j > p_i + \beta z_i + a_{ij} \quad (3.20)$$

$$Q_{ij} = -\sqrt{a_{ij}/b_{ij}} \quad \text{for } p_i + \beta z_i > p_j + \beta z_j \quad (3.21)$$

$$Q_{ij} = -\sqrt{(p_i - p_j + a_{ij} + \beta(z_i - z_j))/b_{ij}} \quad \text{otherwise} \quad (3.22)$$

Here:

$$\beta = \rho g \quad (3.23)$$

3.3.2 For Gas

The flow rate across the compressor from node j to node i is:

$$Q_{sc-ji} = \left(\frac{1}{\psi_{sc}} \right) \frac{W_{c-ji}}{\left\{ \frac{k}{k-1} \zeta T_j \left[\left(\frac{p_i}{p_j} \right)^{(k-1)/k} - 1 \right] + \omega_{ji} \right\}} \quad (3.24)$$

The flow rate across the compressor from node i to node j is:

$$Q_{sc-ij} = \left(\frac{1}{\psi_{sc}} \right) \frac{-W_{c-ij}}{\left\{ \frac{k}{k-1} \zeta T_i \left[\left(\frac{p_j}{p_i} \right)^{(k-1)/k} - 1 \right] + \omega_{ij} \right\}} \quad (3.25)$$

Here:

$$\psi_{sc} = \frac{p_{sc} M}{Z_{sc} R T_{sc}} \quad (3.26)$$

$$\zeta = \frac{Z_{avg} R}{M} \quad (3.27)$$

$$\omega_{ji} = g(z_i - z_j) \quad (3.28)$$

$$\omega_{ij} = g(z_j - z_i) \quad (3.29)$$

$$k = \frac{C_P}{C_V} \quad (3.30)$$

3.4 Pipeline Flow with Partial Derivatives

In the followings relatively small variations of the Fanning friction factor are ignored.

The partial derivatives for the Newton-Raphson method with respect to p_j and p_i are given as follows:

3.4.1 For Liquid

Table 3.1 Liquid flow rate with partial derivatives

$y = p_j - p_i + \beta(z_j - z_i)$	$Q_{ji} = \sqrt{y / (\alpha_{ji} f_{F_{ji}})}$ $y > 0$	$Q_{ij} = -\sqrt{-y / (\alpha_{ji} f_{F_{ji}})}$ $y < 0$
$\frac{\partial Q}{\partial p_j}$	$\frac{\partial Q_{ji}}{\partial p_j} = 0.5 \sqrt{1 / (\alpha_{ji} f_{F_{ji}} y)}$	$\frac{\partial Q_{ij}}{\partial p_j} = 0.5 \sqrt{-1 / (\alpha_{ji} f_{F_{ji}} y)}$
$\frac{\partial Q}{\partial p_i}$	$\frac{\partial Q_{ji}}{\partial p_i} = -0.5 \sqrt{1 / (\alpha_{ji} f_{F_{ji}} y)}$	$\frac{\partial Q_{ij}}{\partial p_i} = -0.5 \sqrt{-1 / (\alpha_{ji} f_{F_{ji}} y)}$

3.4.2 For Gas

Inclined Flow ($z_i \neq z_j$):

Table 3.2 Inclined flow rate with partial derivatives

$Q_{sc-ji} = \frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{w}{\delta_{ji}(\phi_{ji} - 1)}}$ <p>In which: $w = p_j^2 - \phi_{ji} p_i^2$,</p> <p>and $w > 0$</p>	$Q_{sc-ij} = -\frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{-w}{\delta_{ji}(\phi_{ji} - 1)}}$ <p>In which: $w = p_j^2 - \phi_{ji} p_i^2$,</p> <p>and $w < 0$</p>
$\frac{\partial Q_{sc-ji}}{\partial p_j} = \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji}(\phi_{ji} - 1)} \frac{p_j}{\sqrt{\frac{w}{\delta_{ji}(\phi_{ji} - 1)}}$	$\frac{\partial Q_{sc-ij}}{\partial p_j} = \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji}(\phi_{ji} - 1)} \frac{p_j}{\sqrt{\frac{-w}{\delta_{ji}(\phi_{ji} - 1)}}$
$\frac{\partial Q_{sc-ji}}{\partial p_i} = \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji}(\phi_{ji} - 1)} \frac{-\phi_{ji} p_i}{\sqrt{\frac{w}{\delta_{ji}(\phi_{ji} - 1)}}$	$\frac{\partial Q_{sc-ij}}{\partial p_i} = \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji}(\phi_{ji} - 1)} \frac{-\phi_{ji} p_i}{\sqrt{\frac{-w}{\delta_{ji}(\phi_{ji} - 1)}}$

Horizontal Flow ($z_i = z_j$):

Table 3.3 Horizontal flow rate with partial derivatives

$Q_{sc-ji} = \frac{A_{ji}}{\psi_{sc}} \sqrt{\xi_{ji} (p_j^2 - p_i^2)}$ <p>In Case: $p_j > p_i$</p>	$Q_{sc-ij} = -\frac{A_{ji}}{\psi_{sc}} \sqrt{-\xi_{ji} (p_j^2 - p_i^2)}$ <p>In Case: $p_j < p_i$</p>
$\frac{\partial Q_{sc-ji}}{\partial p_j} = \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{\xi_{ji} p_j}{\sqrt{\xi_{ji} (p_j^2 - p_i^2)}}$	$\frac{\partial Q_{sc-ij}}{\partial p_j} = \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{\xi_{ji} p_j}{\sqrt{-\xi_{ji} (p_j^2 - p_i^2)}}$
$\frac{\partial Q_{sc-ji}}{\partial p_i} = \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{-\xi_{ji} p_i}{\sqrt{\xi_{ji} (p_j^2 - p_i^2)}}$	$\frac{\partial Q_{sc-ij}}{\partial p_i} = \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{-\xi_{ji} p_i}{\sqrt{-\xi_{ji} (p_j^2 - p_i^2)}}$

3.5 Equipment Flow with Partial Derivatives

The partial derivatives of non-zero Q_{ji} and Q_{ij} for the Newton-Raphson method with respect to p_j and p_i respectively are given as follows:

3.5.1 For Liquid

Table 3.4 Non-zero liquid flow rate across a pump with partial derivatives

$Q_{ji} = \sqrt{w_{ji}/b_{ji}}$ $w_{ji} = p_j - p_i + a_{ji} + \beta(z_j - z_i)$ <p>In which: $C_{ji} = 2$</p>	$Q_{ij} = -\sqrt{w_{ij}/b_{ij}}$ $w_{ij} = p_i - p_j + a_{ij} + \beta(z_i - z_j)$ <p>In which: $C_{ij} = 2$</p>
$\frac{\partial Q_{ji}}{\partial p_j} = \frac{1}{2} \sqrt{1/(b_{ji} w_{ji})}$	$\frac{\partial Q_{ij}}{\partial p_j} = \frac{1}{2} \sqrt{1/(b_{ij} w_{ij})}$
$\frac{\partial Q_{ji}}{\partial p_i} = -\frac{1}{2} \sqrt{1/(b_{ji} w_{ji})}$	$\frac{\partial Q_{ij}}{\partial p_i} = -\frac{1}{2} \sqrt{1/(b_{ij} w_{ij})}$

3.5.2 For Gas

Table 3.5 Gas flow rate across a compressor with partial derivatives

$Q_{sc-j_i} = \left(\frac{1}{\psi_{sc}} \right) \frac{W_{c-j_i}}{CR_{j_i}}$ $CR_{j_i} = \frac{k}{k-1} \zeta T_j \left[\left(\frac{p_i}{p_j} \right)^{(k-1)/k} - 1 \right] + \omega_{j_i}$ <p style="text-align: center;">In which: $C_{j_i} = 3$</p>	$Q_{sc-i_j} = \left(\frac{1}{\psi_{sc}} \right) \frac{-W_{c-i_j}}{CR_{i_j}}$ $CR_{i_j} = \frac{k}{k-1} \zeta T_i \left[\left(\frac{p_j}{p_i} \right)^{(k-1)/k} - 1 \right] + \omega_{i_j}$ <p style="text-align: center;">In which: $C_{i_j} = 3$</p>
$\frac{\partial Q_{sc-j_i}}{\partial p_j} = \frac{T_j \zeta}{p_j \psi_{sc}} \frac{W_{c-j_i}}{CR_{j_i}^2} \left(\frac{p_i}{p_j} \right)^{(k-1)/k}$	$\frac{\partial Q_{sc-i_j}}{\partial p_j} = \frac{T_i \zeta}{p_j \psi_{sc}} \frac{W_{c-i_j}}{CR_{i_j}^2} \left(\frac{p_j}{p_i} \right)^{(k-1)/k}$
$\frac{\partial Q_{sc-j_i}}{\partial p_i} = \frac{T_j \zeta}{p_i \psi_{sc}} \frac{-W_{c-j_i}}{CR_{j_i}^2} \left(\frac{p_i}{p_j} \right)^{(k-1)/k}$	$\frac{\partial Q_{sc-i_j}}{\partial p_i} = \frac{T_i \zeta}{p_i \psi_{sc}} \frac{-W_{c-i_j}}{CR_{i_j}^2} \left(\frac{p_j}{p_i} \right)^{(k-1)/k}$

3.6 Conversion Units

Table 3.6 British and SI units

Quantity	British units	SI units
p, p_{sc}	psig	bar
Q	gpm	m^3/hr
Q_{sc}	MMscfd	MMscmd
ρ	lb_m/ft^3	kg/m^3
μ	centipoise	mPa – s
L	ft	m
D	inch	mm
z	ft	m
ε	ft	mm
g	ft/sec^2	m/sec^2
T, T_{sc}	$^{\circ}F$	$^{\circ}C$
M	lb_m	kg
R	$ft\ lb_f/lb\ mole\ ^{\circ}R$	$J/kmole\ ^{\circ}K$
Z_{avg}, Z_{sc}	none	none
f_F	none	none
Re	none	none

3.6.1 For Liquid

Table 3.7 Conversion units for α_{ji} , β and Re_{ji}

British units	SI units
$\alpha_{ji} = \frac{32 * (12)^5 * \rho L_{ji}}{\pi^2 * 144 * 32.2 * (7.48 * 60)^2 * D_{ji}^5}$	$\alpha_{ji} = \frac{32 * 10^{15} * \rho L_{ji}}{\pi^2 * (3600)^2 * 1.01325 * 10^5 * D_{ji}^5}$
$\beta = \frac{\rho}{144}$	$\beta = \frac{9.81 * \rho}{1.01325 * 10^5}$
$Re_{ji} = \frac{4 * 12 * 10^5 * \rho Q_{ji}}{7.48 * 60 * 32.2 * 2.089 * \pi \mu D_{ji}}$	$Re_{ji} = \frac{4 * 10^6 * \rho Q_{ji}}{3600 * \pi \mu D_{ji}}$

3.6.2 For Gas

Inclined Flow ($z_i \neq z_j$)

Table 3.8 Conversion units for λ_{ji} , δ_{ji} and ϕ_{ji}

British units	SI units
$\lambda_{ji} = \left(\frac{24 * 60 * 60}{(10)^6 * (12)^2 * Z_{avg} T} \right) \left(\frac{\pi D_{ji}^2}{4} \right)$	$\lambda_{ji} = \left(\frac{24 * 3600}{(10)^6 * (10)^6 * Z_{avg} T} \right) \left(\frac{\pi D_{ji}^2}{4} \right)$
$\delta_{ji} = \left(\frac{2 * f_{F_{ji}} L_{ji}}{(32.2 * 12)^3 * D_{ji} (z_i - z_j)} \right)$	$\delta_{ji} = \left(\frac{2 * 10^3 * f_{F_{ji}} L_{ji}}{9.81 * (1.01325 * 10^5)^2 * D_{ji} (z_i - z_j)} \right)$
$\phi_{ji} = \exp \left(\frac{2 * M (z_i - z_j)}{1545.3 * Z_{avg} T} \right)$	$\phi_{ji} = \exp \left(\frac{2 * 9.81 * M (z_i - z_j)}{8314.3 * Z_{avg} T} \right)$
$\psi_{sc} = \frac{32.2 * (12)^2 p_{sc}}{Z_{sc} T_{sc}}$	$\psi_{sc} = \frac{1.01325 * 10^5 * p_{sc}}{Z_{sc} T_{sc}}$

Horizontal Flow ($z_i = z_j$)

Table 3.9 Conversion units for A_{ji} , ξ_{ji} and ψ_{sc}

British units	SI units
$A_{ji} = \frac{24 * 60 * 60 * \pi D_{ji}^2}{(10)^6 * 4 * 144}$	$A_{ji} = \frac{24 * 3600 * \pi D_{ji}^2}{(10)^6 * 4 * 10^6}$
$\xi_{ji} = \left(\frac{32.2 * (12)^3 * MD_{ji}}{4 * 1545.3 * Z_{avg} Tf_{F_{ji}} L_{ji}} \right)$	$\xi_{ji} = \left(\frac{(1.01325 * 10^5)^2 MD_{ji}}{4 * 8314.3 * 10^3 * Z_{avg} Tf_{F_{ji}} L_{ji}} \right)$
$\psi_{sc} = \frac{(12)^2 * p_{sc} M}{1545.3 * Z_{sc} T_{sc}}$	$\psi_{sc} = \frac{1.01325 * 10^5 * p_{sc} M}{8314.3 * Z_{sc} T_{sc}}$

Note:

$$^{\circ}R = ^{\circ}F + 459.67,$$

$$^{\circ}K = ^{\circ}C + 273.15$$

$$p_{absolute} = p_{gauge} + 14.73,$$

$$p_{avg-ji} = \frac{2}{3} \left(\frac{p_j^3 - p_i^3}{p_j^2 - p_i^2} \right)$$

British unit:

$$Re_{ji} = \left(\frac{4 * 12 * 10^5 * (12)^2 * (10)^6 * Q_{sc-ji}}{32.2 * 2.089 * 1545.3 * 24 * 3600 * \pi \mu D_{ji}} \right) \left(\frac{p_{avg-ji} M}{Z_{avg} T} \right)$$

SI unit:

$$Re_{ji} = \left(\frac{4 * 10^6 * 1.01325 * (10)^5 * (10)^6 * Q_{sc-ji}}{8314.3 * 24 * 3600 * \pi \mu D_{ji}} \right) \left(\frac{p_{avg-ji} M}{Z_{avg} T} \right)$$

3.7 Nodal Material Balance Equations

The nodal material balance equations for all nodes i at which the pressure p_i , is not specified (for $T_i \neq 1$) can be described as follows:

For steady-state, the sum of the flows into any node i must be zero. That is:

$$F_i(\mathbf{P}) = 0, \quad (3.31)$$

Here:

$F_i(\mathbf{P})$ is the net flow into any node i .

$$\mathbf{P} = [p_1, p_2, \dots, p_n]^T$$

3.7.1 For Liquid

$$\begin{aligned} F_i(\mathbf{P}) = & \text{injection rate (or withdrawal rate)} & (3.32) \\ & + \text{net flow in from neighboring nodes to } i \text{ by pipeline} \\ & + \text{net flow in from neighboring nodes to } i \text{ from pumps} \\ & - \text{net flow out to neighboring nodes from } i \text{ through pumps} \\ = & 0 \end{aligned}$$

The equation for $F_i(\mathbf{P})$ becomes:

$$\begin{aligned} F_i(\mathbf{P}) = & V_i (\text{injection (positive) or withdrawal (negative) rate}) \\ & + \sum_{j, C_{ji}=1} Q_{ji} (\text{two pipeline cases}) \\ & + \sum_{j, C_{ji}=2} Q_{ji} (\text{three cases for pumping in}) \\ & + \sum_{j, C_{ij}=2} Q_{ji} (\text{three cases for pumping out}) \end{aligned} \quad (3.33)$$

3.7.2 For Gas

$$\begin{aligned}
 F_i(\mathbf{P}) = & \text{injection rate (or withdrawal rate)} & (3.34) \\
 & + \text{net flow in from neighboring nodes to } i \text{ by pipeline} \\
 & + \text{net flow in from neighboring nodes to } i \text{ from compressors} \\
 & - \text{net flow out to neighboring nodes from } i \text{ through compressors} \\
 = & 0
 \end{aligned}$$

The equation for $F_i(\mathbf{P})$ becomes:

$$\begin{aligned}
 F_i(\mathbf{P}) = & V_i (\text{injection (positive) or withdrawal (negative) rate}) \\
 & + \sum_{j, C_{ji}=1} Q_{ji} (\text{two pipeline cases for inclined flow}) \\
 & + \sum_{j, C_{ji}=1} Q_{ji} (\text{two pipeline cases for horizontal flow}) \\
 & + \sum_{j, C_{ji}=3} Q_{ji} (\text{the case for compressor coming in}) \\
 & + \sum_{j, C_{ij}=3} Q_{ji} (\text{the case for compressor going out})
 \end{aligned} \tag{3.35}$$

3.8 Newton-Raphson Method

The simultaneous nonlinear equations in the unknown pressures that are obtained from nodal material balances at all nodes i are solved by the iterative Newton-Raphson method as follows:

1. Suppose we have an initial estimate of $f_{F_{ji}}$ for all connections between node j and node i such that $C_{ji} = 1$.

2. Suppose we also know the approximate and specified pressure p_i , at all nodes i .

3. The next step is to find the appropriate partial derivatives of the functions $F_i(\mathbf{P})$, ($i = 1, 2, \dots, n$) with respect to p_j , ($j = 1, 2, \dots, n$) which are then stored as the elements of the left hand side coefficient matrix, Φ of the simultaneous linear equations:

$$\Phi(\mathbf{P})\delta\mathbf{P} = -\mathbf{F}(\mathbf{P}) \quad (3.36)$$

In Eqn. (3.36), the right hand side vector is defined as:

$$\mathbf{F}(\mathbf{P}) = [F_1(\mathbf{P}), F_2(\mathbf{P}), F_3(\mathbf{P}), \dots, F_n(\mathbf{P})]^t \quad (3.37)$$

where the correction vector $\delta\mathbf{P}$ is the solution of the simultaneous linear equations and a representative element of the coefficient matrix is:

$$\Phi(\mathbf{P}) = F_{ij}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_j}, \quad 1 \leq i, j \leq n \quad (3.38)$$

3.8.1 For Liquid

The partial derivative of $F_i(\mathbf{P})$ with respect to p_j is given by one of following forms:

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{ji}}{\partial p_j} = \left\{ \begin{array}{ll} 0.5\sqrt{1/(\alpha_{ji} f_{F_{ji}} y)} & \text{for } y > 0 \\ 0.5\sqrt{-1/(\alpha_{ji} f_{F_{ji}} y)} & \text{for } y < 0 \end{array} \right\} \text{ if } C_{ji} = 1$$

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{ji}}{\partial p_j} = \frac{1}{2} \sqrt{1/(b_{ji} w_{ji})} \quad \text{if } C_{ji} = 2$$

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{ji}}{\partial p_j} = \frac{1}{2} \sqrt{1/(b_{ij} w_{ij})} \quad \text{if } C_{ij} = 2$$

Here ($i \neq j$):

$$F_{ij}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_j}, \quad 1 \leq i, j \leq n \quad (3.39)$$

The partial derivatives of $F_i(\mathbf{P})$ with respect to p_i are given by summation as follows:

$$\begin{aligned} F_{ii}(\mathbf{P}) &= \sum_{j, C_{ji}=1} \frac{\partial Q_{ji}}{\partial p_i} \text{ (two pipeline cases)} \\ &+ \sum_{j, C_{ji}=2} \frac{\partial Q_{ji}}{\partial p_i} \text{ (three cases for pumping in)} \\ &+ \sum_{j, C_{ij}=2} \frac{\partial Q_{ji}}{\partial p_i} \text{ (three cases for pumping out)} \end{aligned} \quad (3.40)$$

$$\begin{aligned} F_{ii}(\mathbf{P}) &= \left\{ \begin{array}{ll} \sum_{j, C_{ji}=1} -0.5 \sqrt{1/(\alpha_{ji} f_{F_{ji}} y)} & \text{for } y > 0 \\ \sum_{j, C_{ji}=1} -0.5 \sqrt{-1/(\alpha_{ji} f_{F_{ji}} y)} & \text{for } y < 0 \end{array} \right\} \\ &+ \sum_{j, C_{ji}=2} -\frac{1}{2} \sqrt{1/(b_{ji} w_{ji})} \\ &+ \sum_{j, C_{ij}=2} -\frac{1}{2} \sqrt{1/(b_{ij} w_{ij})} \end{aligned} \quad (3.41)$$

Here:

$$F_{ii}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_i}, \quad 1 \leq i \leq n \quad (3.42)$$

3.8.2 For Gas

The partial derivative of $F_i(\mathbf{P})$ with respect to p_j is given by one of following forms:

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{sc-ji}}{\partial p_j} = \left\{ \begin{array}{l} \left. \begin{array}{l} \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji} (\phi_{ji} - 1)} \frac{p_j}{\sqrt{\frac{w}{\delta_{ji} (\phi_{ji} - 1)}}}} \quad \text{for } w > 0 \\ \frac{\lambda_{ji}}{\psi_{sc} \delta_{ji} (\phi_{ji} - 1)} \frac{p_j}{\sqrt{\frac{-w}{\delta_{ji} (\phi_{ji} - 1)}}}} \quad \text{for } w < 0 \end{array} \right\} \quad \text{if } C_{ji} = 1 \\ \left. \begin{array}{l} \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{\xi_{ji} p_j}{\sqrt{\xi_{ji} (p_j^2 - p_i^2)}} \quad \text{for } p_j > p_i \\ \left(\frac{A_{ji}}{\psi_{sc}} \right) \frac{\xi_{ji} p_j}{\sqrt{-\xi_{ji} (p_j^2 - p_i^2)}} \quad \text{for } p_j < p_i \end{array} \right\} \end{array} \right.$$

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{sc-ji}}{\partial p_j} = \frac{T_j \zeta}{p_j \psi_{sc}} \frac{W_{c-ji}}{\left\{ \frac{k}{k-1} \zeta T_j \left[\left(\frac{p_i}{p_j} \right)^{(k-1)/k} - 1 \right] + \omega_{ji} \right\}^2} \left(\frac{p_i}{p_j} \right)^{(k-1)/k} \quad \text{if } C_{ji} = 3$$

$$F_{ij}(\mathbf{P}) = \frac{\partial Q_{sc-ji}}{\partial p_j} = \frac{T_i \zeta}{p_j \psi_{sc}} \frac{W_{c-ij}}{\left\{ \frac{k}{k-1} \zeta T_i \left[\left(\frac{p_j}{p_i} \right)^{(k-1)/k} - 1 \right] + \omega_{ij} \right\}^2} \left(\frac{p_j}{p_i} \right)^{(k-1)/k} \quad \text{if } C_{ij} = 3$$

The partial derivatives of $F_i(\mathbf{P})$ with respect to p_i are given by summation as follows:

$$\begin{aligned} F_{ii}(\mathbf{P}) &= \sum_{j, C_{ji}=1} \frac{\partial Q_{ji}}{\partial p_i} \quad (\text{two pipeline cases, for both inclined and horizontal flow}) \\ &+ \sum_{j, C_{ji}=3} \frac{\partial Q_{ji}}{\partial p_i} \quad (\text{the case for compressor coming in}) \\ &+ \sum_{j, C_{ij}=3} \frac{\partial Q_{ji}}{\partial p_i} \quad (\text{the case for compressor going out}) \end{aligned}$$

$$\begin{aligned}
F_{ii}(\mathbf{P}) = \sum_{j, C_{ji}=1} & \left\{ \begin{array}{l} z_i \neq z_j \left\{ \begin{array}{l} \frac{\lambda_{ji}}{\Psi_{sc} \delta_{ji} (\phi_{ji} - 1)} \frac{-\phi_{ji} p_i}{\sqrt{\frac{w}{\delta_{ji} (\phi_{ji} - 1)}}} \quad \text{for } w > 0 \\ \frac{\lambda_{ji}}{\Psi_{sc} \delta_{ji} (\phi_{ji} - 1)} \frac{-\phi_{ji} p_i}{\sqrt{\frac{-w}{\delta_{ji} (\phi_{ji} - 1)}}} \quad \text{for } w < 0 \end{array} \right. \\ z_i = z_j \left\{ \begin{array}{l} \frac{A_{ji}}{\Psi_{sc} \sqrt{\xi_{ji} (p_j^2 - p_i^2)}} \frac{-\xi_{ji} p_i}{\sqrt{\xi_{ji} (p_j^2 - p_i^2)}} \quad \text{for } p_j > p_i \\ \frac{A_{ji}}{\Psi_{sc} \sqrt{-\xi_{ji} (p_j^2 - p_i^2)}} \frac{-\xi_{ji} p_i}{\sqrt{-\xi_{ji} (p_j^2 - p_i^2)}} \quad \text{for } p_j < p_i \end{array} \right. \end{array} \right. \\
+ \sum_{j, C_{ji}=3} \frac{T_j}{p_i} \frac{\zeta}{\Psi_{sc}} \frac{-W_{c-ji}}{\left\{ \frac{k}{k-1} \zeta T_j \left[\left(\frac{p_i}{p_j} \right)^{(k-1)/k} - 1 \right] + \omega_{ji} \right\}^2} \left(\frac{p_i}{p_j} \right)^{(k-1)/k} \\
+ \sum_{j, C_{ji}=3} \frac{T_i}{p_i} \frac{\zeta}{\Psi_{sc}} \frac{-W_{c-ij}}{\left\{ \frac{k}{k-1} \zeta T_i \left[\left(\frac{p_j}{p_i} \right)^{(k-1)/k} - 1 \right] + \omega_{ij} \right\}^2} \left(\frac{p_j}{p_i} \right)^{(k-1)/k}
\end{aligned}$$

4. Use LU decomposition of the Gaussian elimination method with column pivoting only to solve the simultaneous linear equation with $\Phi(\mathbf{P})$ as the left hand side coefficient matrix.

5. Back substitution to find out the correction vector $\delta\mathbf{P}$. Also using a mathematical technique to improve the stability of the method at all nodes i by factor σ_i , in the correction as follows:

$$\delta p_i = \delta p_i^* \sigma_i \quad (3.43)$$

Where:

δp_i is the value of the correction actually applied.

δp_i^* is the value of the correction computed from the Newton-Raphson method.

It is recommended that $\sigma_i = 0.5$ is the best value to use in order to ensure convergence for the first iteration. In subsequent iterations, the value of " σ_i " is determined as below:

For $A_i \leq -1$	$\sigma_i = 0.5 A_i $
For $-1 < A_i < 0$	$\sigma_i = 0.4 - 0.15 A_i $
For $0 < A_i < 1$	$\sigma_i = 0.4 + 0.15 A_i $
For $A_i \geq 1$	$\sigma_i = 0.5$

Here A_i is computed by using the δp_i for the current and previous iterations as follows:

$$A_i = \frac{\delta p_i^{k+1}}{\delta p_i^k} \quad i = 1, 2, \dots, n \quad (3.44)$$

In which:

δp_i^{k+1} is the correction to p_i for the current iteration.

δp_i^k is the correction of p_i for the previous iteration.

Note:

The user has to do some experimentation to obtain the coefficients of c_1, c_2, c_3, c_4, c_5 and c_6 for the factor σ_i , in accordance with his or her own system. (generally, $0.0 \leq c_1, c_2, c_3, c_4, c_5, c_6 \leq 1.0$)

6. Check for convergence after the corrections δp_i have been made at all nodes i (improved by the factor to avoid instability in item 5. if necessary) according to some criterion such as:

$$|\delta p_i| < \lambda \quad i = 1, 2, \dots, n \quad (3.45)$$

7. If the corrections δP do not satisfy the convergence condition, the current vector of pressures is modified according to:

$$\mathbf{P}_{k+1} = \mathbf{P}_k + \delta \mathbf{P}_k \quad (3.46)$$

Here:

\mathbf{P}_k is the current vector (or set) of pressures.

$\delta \mathbf{P}_{k+1}$ is the updated set of pressures for use the next iteration.

$\delta \mathbf{P}_k$ is the set of pressure corrections just computed.

8. With these new pressures \mathbf{P}_{k+1} , from equation (3.46), the updated flow rates Q_{ji} , can be calculated with the old Fanning friction factor $f_{F_{ji}}^k$, for all pipeline segments as follows:

For Liquid:

$$Q_{ji} = \sqrt{y / (\alpha_{ji} f_{F_{ji}})} \quad \text{for } y > 0 \quad (3.47)$$

$$Q_{ij} = -\sqrt{-y / (\alpha_{ji} f_{F_{ji}})} \quad \text{for } y < 0 \quad (3.48)$$

For Gas:

Inclined Flow ($z_i \neq z_j$):

$$Q_{sc-ji} = \frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{w}{\delta_{ji} (\phi_{ji} - 1)}} \quad \text{for } w > 0 \quad (3.49)$$

$$Q_{sc-ij} = -\frac{\lambda_{ji}}{\psi_{sc}} \sqrt{\frac{-w}{\delta_{ji} (\phi_{ji} - 1)}} \quad \text{for } w < 0 \quad (3.50)$$

Horizontal Flow ($z_i = z_j$):

$$Q_{sc-ji} = \frac{A_{ji}}{\psi_{sc}} \sqrt{\xi_{ji} (p_j^2 - p_i^2)} \quad \text{for } p_j > p_i \quad (3.51)$$

$$Q_{sc-ij} = -\frac{A_{ji}}{\psi_{sc}} \sqrt{-\xi_{ji} (p_j^2 - p_i^2)} \quad \text{for } p_j < p_i \quad (3.52)$$

9. The Reynolds numbers Re_{ji} , are computed for all pipeline segments as follows:

For Liquid:

$$Re_{ji} = \frac{4 * 12 * 10^5 * \rho Q_{ji}}{7.48 * 60 * 32.2 * 2.089 * \pi \mu D_{ji}} \quad \text{for British units}$$

$$Re_{ji} = \frac{4 * 10^6 \rho Q_{ji}}{3600 \pi \mu D_{ji}} \quad \text{for SI units}$$

For Gas:

British units:

$$Re_{ji} = \frac{4 * 12 * 10^5 * (12)^2 * Q_{ji}}{7.48 * 60 * 32.2 * 2.089 * 1545.3 * \pi \mu D_{ji}} \frac{p_{avg-ji} M}{Z_{avg} T}$$

SI units:

$$Re_{ji} = \frac{4 * 10^6 * 1.01325 * 10^5 * Q_{ji}}{3600 * 8314.3 * \pi \mu D_{ji}} \frac{p_{avg-ji} M}{Z_{avg} RT}$$

10. The program updates the Fanning friction factors $f_{F_{ji}}^k$, as functions of the Reynolds number and roughness ratio in all pipeline segments as follows:

For turbulent flow ($Re_{ji} > 4000$):

$$f_{F_{ji}} = \left\{ -1.737 \ln \left[0.269 \frac{\epsilon_{ji}}{D_{ji}} - \frac{2.185}{Re_{ji}} \ln \left(0.269 \frac{\epsilon_{ji}}{D_{ji}} + \frac{14.5}{Re_{ji}} \right) \right] \right\}^{-2} \quad (3.53)$$

For laminar flow ($Re_{ji} \leq 2000$):

$$f_{F_{ji}} = \frac{16}{Re_{ji}} \quad (3.54)$$

11. The sequence of calculations given above is repeated for successive iterations $k = 1, 2, 3 \dots$ until convergence occurs according to some predetermined criterion such as:

$$\left| (\mathbf{P}_{k+1})_i - (\mathbf{P}_k)_i \right| < \lambda \quad \text{for all } i = 1, 2, \dots, n \quad (3.55)$$

or until a specified maximum number of iterations k_{\max} , has been exceeded.

12. If the i th node type is $T_i = 1$ (pressure specified), it can be included in the Newton-Raphson method by using it as an unknown in the simultaneous linear equations and setting its correction δP , to zero. This is achieved by setting:

$$\left. \begin{array}{l} F_i(\mathbf{P}) = 0 \\ F_{ii}(\mathbf{P}) = 1 \\ F_{ij}(\mathbf{P}) = 0 \end{array} \right\} \text{For } T_i = 1 \quad (3.56)$$

3.9 Terminal Node with Specified Injection Rate

Consider the special case of a terminal node i with a specified injection rate V_i , as follows:

For Liquid:

The net flow into terminal node i must equal zero, so that:

$$Q_{ji} = -V_i \quad (3.57)$$

In the case of pipeline connection, the flow rates from Eqns. (3.1) and (3.2) can be represented by one equation instead of two as follows:

$$y = \pm \alpha_{ji} f_{F_{ji}} Q_{ji}^2 \quad (3.58)$$

$$y = -\alpha_{ji} f_{F_{ji}} V_i |V_i| \quad (3.59)$$

Define:

$$F_i(\mathbf{P}) = -\alpha_{ji} f_{F_{ji}} V_i |V_i| - y = 0 \quad (3.60)$$

Therefore:

$$-F_i(\mathbf{P}) = \alpha_{ji} f_{F_{ji}} V_i |V_i| + y \quad (3.61)$$

The partial derivatives of the function $F_i(\mathbf{P})$ with respect to p_j and p_i are given by:

$$F_{ij}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_j} = -1 \quad (3.62)$$

$$F_{ii}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_i} = 1 \quad (3.63)$$

For Gas:

The net flow into terminal node i at standard conditions must equal zero, thus:

$$Q_{sc-ji} = -V_i \quad (3.64)$$

Inclined Flow ($z_i \neq z_j$):

For pipeline connection, the flow rates from Eqns. (3.5) and (3.6) can be placed by one equation as follows:

$$Q_{sc-ji}^2 = \pm \frac{\lambda_{ji}^2}{\psi_{sc}^2 \delta_{ji}} \frac{w}{(\phi_{ji} - 1)} \quad (3.65)$$

$$-V_i |V_i| = \left(\frac{\lambda_{ji}}{\psi_{sc}} \right)^2 \frac{w}{\delta_{ji} (\phi_{ji} - 1)} \quad (3.66)$$

Rearrangement gives:

$$-\psi_{sc}^2 V_i |V_i| = \frac{\lambda_{ji}^2 w}{\delta_{ji} (\phi_{ji} - 1)} \quad (3.67)$$

Define:

$$F_i(\mathbf{P}) = -\psi_{sc}^2 V_i |V_i| - \frac{\lambda_{ji}^2 w}{\delta_{ji} (\phi_{ji} - 1)} = 0 \quad (3.68)$$

Therefore:

$$-F_i(\mathbf{P}) = \psi_{sc}^2 V_i |V_i| + \frac{\lambda_{ji}^2 w}{\delta_{ji} (\phi_{ji} - 1)} \quad (3.69)$$

The partial derivatives of the function $F_i(\mathbf{P})$ with respect to p_j and p_i are given by:

$$F_{ij}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_j} = \frac{-2\lambda_{ji}^2 p_j}{\delta_{ji} (\phi_{ji} - 1)} \quad (3.70)$$

$$F_{ii}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_i} = \frac{2\lambda_{ji}^2 \phi_{ji} p_i}{\delta_{ji} (\phi_{ji} - 1)} \quad (3.71)$$

Horizontal Flow ($z_i = z_j$):

In the same manner of inclined flow, the flow rates from Eqns. (3.12) and (3.13) can be reduced to one equation as follows:

$$Q_{sc-ji} = \pm \frac{A_{ji}^2}{\psi_{sc}^2} \xi_{ji} (p_j^2 - p_i^2) \quad (3.72)$$

$$-V_i |V_i| = \left(\frac{A_{ji}}{\psi_{sc}} \right)^2 \xi_{ji} (p_j^2 - p_i^2) \quad (3.73)$$

Rearrangement gives:

$$-\psi_{sc}^2 V_i |V_i| = A_{ji}^2 \xi_{ji} (p_j^2 - p_i^2) \quad (3.74)$$

Define:

$$F_i(\mathbf{P}) = -\psi_{sc}^2 V_i |V_i| - A_{ji}^2 \xi_{ji} (p_j^2 - p_i^2) = 0 \quad (3.75)$$

Thus:

$$-F_i(\mathbf{P}) = \psi_{sc}^2 V_i |V_i| + A_{ji}^2 \xi_{ji} (p_j^2 - p_i^2) \quad (3.76)$$

The partial derivatives of the function $F_i(\mathbf{P})$ with respect to p_j and p_i are given by:

$$F_{ij}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_j} = -2p_j \xi_{ji} A_{ji}^2 \quad (3.77)$$

$$F_{ii}(\mathbf{P}) = \frac{\partial F_i(\mathbf{P})}{\partial p_i} = 2p_i \xi_{ji} A_{ji}^2 \quad (3.78)$$

3.10 FORTRAN Language

A FORTRAN program (Power Station Version 1.0) is written to accept the above information concerning any network of nodes i and use the Newton-Raphson iterative technique to compute the unknown nodal pressures at all nodes i of node type $T_i = 0$ or $T_i = 3$. The output displays a set of matrices containing the internodal flow rates Q_{ij} , nodal pressures and the Fanning friction factors $f_{F_{ji}}^k$, for all pipeline segments.

3.11 Program Description

A general flow diagram of the program is shown in Fig. 3.1 Subroutine SGEM is used to solve the simultaneous linear equations generated at each new iteration of Newton-Raphson method. Subroutine UP is implemented to generate the next estimates of fanning friction factor after no convergence test.

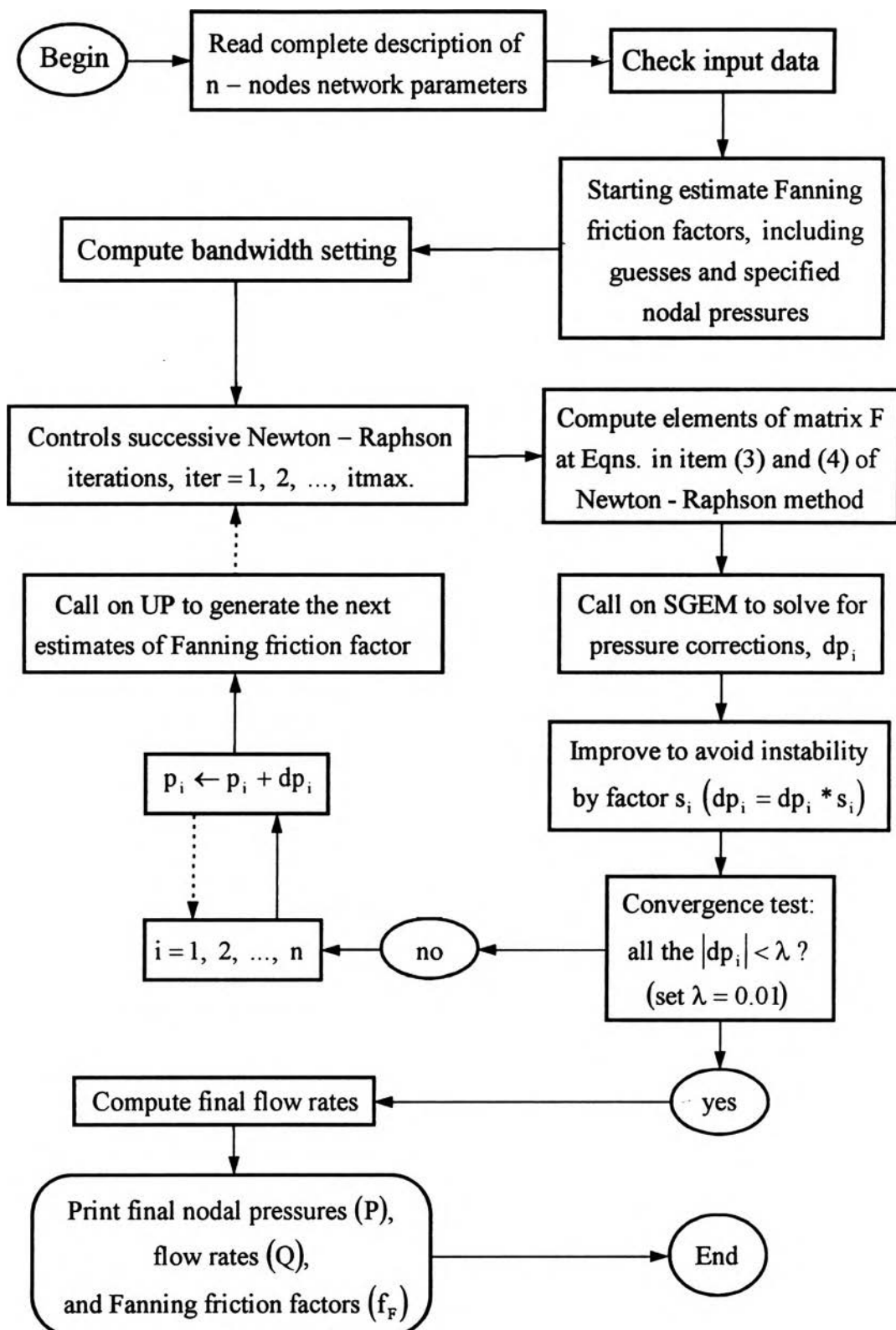


Fig. 3.1 A general flow diagram for fluid network analysis program.