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## **APPENDICES**

## APPENDIX A

The Matrices  $\mathbf{R}(\xi, z, s)$  and  $\mathbf{S}(\xi, z, s)$  in equations (3.15) and (3.16) respectively, are given by

$$\mathbf{R}(\xi, z, s) = \begin{bmatrix} -\xi\delta e^{\gamma z} & -\xi\delta e^{-\gamma z} & a_1 z e^{\xi z} & -a_1 z e^{-\xi z} & e^{\xi z} & e^{-\xi z} \\ \gamma\delta e^{\gamma z} & -\gamma\delta e^{-\gamma z} & -\left(a_1 z - \frac{a_2}{\xi}\right) e^{\xi z} & -\left(a_1 z + \frac{a_2}{\xi}\right) e^{-\xi z} & -e^{\xi z} & e^{-\xi z} \\ 2\mu a_3 \eta e^{\gamma z} & 2\mu a_3 \eta e^{-\gamma z} & -2\mu a_4 \eta e^{\xi z} & -2\mu a_4 \eta e^{-\xi z} & 0 & 0 \end{bmatrix} \quad (\text{A-1})$$

$$\mathbf{S}(\xi, z, s) = 2\mu \begin{bmatrix} -\gamma\xi\delta e^{\gamma z} & \gamma\xi\delta e^{-\gamma z} & \left(a_1\xi z - \frac{1}{2}\right) e^{\xi z} & \left(a_1\xi z + \frac{1}{2}\right) e^{-\xi z} & \xi e^{\xi z} & -\xi e^{-\xi z} \\ \xi^2\delta e^{\gamma z} & \xi^2\delta e^{-\gamma z} & (a_4 - a_1\xi z) e^{\xi z} & (a_4 + a_1\xi z) e^{-\xi z} & -\xi e^{\xi z} & -\xi e^{-\xi z} \\ -a_3 K \gamma \frac{\delta}{c} e^{\gamma z} & a_3 K \gamma \frac{\delta}{c} e^{-\gamma z} & a_4 K \xi \frac{\delta}{c} e^{\xi z} & -a_4 K \xi \frac{\delta}{c} e^{-\xi z} & 0 & 0 \end{bmatrix} \quad (\text{A-2})$$

$$\text{where } a_1 = \frac{1}{2(1-2\nu_u)} \quad (\text{A-3})$$

$$a_2 = \frac{(3-4\nu_u)}{2(1-2\nu_u)} \quad (\text{A-4})$$

$$a_3 = \frac{B(1+\nu_u)(1-\nu)}{3(\nu_u - \nu)} \quad (\text{A-5})$$

$$a_4 = \frac{(1-\nu_u)}{(1-2\nu_u)} \quad (\text{A-6})$$

$$c = 2\mu a_3 \kappa \eta \quad (\text{A-7})$$

$$\eta = \frac{B(1+\nu_u)}{3(1-\nu_u)} \quad (\text{A-8})$$

$$\delta = \frac{\eta}{s/c} \quad (\text{A-9})$$

$$\gamma = \sqrt{\xi^2 + \frac{s}{c}} \quad (\text{A-10})$$

## APPENDIX B

This appendix presents the derivation of the strain energy of fictitious elastic pile  $i^{th}$  in Laplace domain, as shown in Figure 2(b). The axial strain corresponding to equation (3.38) can be expressed as

$$\bar{\varepsilon}^i(z, s) = \sum_{m=1}^{Nt} \frac{-(m-1)\bar{\alpha}_m^i(s)}{L} e^{-(m-1)z/L} \quad (\text{B-1})$$

In view of the conventional constitutive relation the strain energy of fictitious elastic pile  $i^{th}$  can be written as

$$Up^i = \frac{1}{2} \int_L^A \bar{\sigma}^i \bar{\varepsilon}^i dA dz = \frac{\pi(a^i)^2}{2} \int_0^L E^i * (\bar{\varepsilon}^i)^2 dz \quad (\text{B-2})$$

Note that since the medium is multilayered, equation (B-3) rewritten in view of equation (B-2) as

$$\begin{aligned} Up_g &= \frac{\pi(a^i)^2}{2} \int_0^{\Delta t_1^i} E_1^i * \sum_{m=1}^{Nt} \sum_{n=1}^{Nt} (m-1)(n-1) \bar{\alpha}_m^i(s) \bar{\alpha}_n^i(s) e^{-(m+n-2)(z/L)} dz \\ &+ \frac{\pi(a^i)^2}{2} \int_{\Delta t_1^i}^{\Delta t_1^i + \Delta t_2^i} E_2^i * \sum_{m=1}^{Nt} \sum_{n=1}^{Nt} (m-1)(n-1) \bar{\alpha}_m^j(s) \bar{\alpha}_n^j(s) e^{-(m+n-2)(z/L)} dz \\ &\vdots \\ &+ \frac{\pi(a^i)^2}{2} \int_{\Delta t_1^i + \Delta t_2^i + \dots + \Delta t_{Ne-1}^i}^{\Delta t_1^i + \Delta t_2^i + \dots + \Delta t_{Ne}^i} E_{Ne}^i * \sum_{m=1}^{Nt} \sum_{n=1}^{Nt} (m-1)(n-1) \bar{\alpha}_m^i(s) \bar{\alpha}_n^i(s) e^{-(m+n-2)(z/L)} dz \end{aligned} \quad (\text{B-3})$$

By integrating equation (B-3), the strain energy of fictitious elastic pile  $i^{th}$  can be obtain in the following form

$$Up^i = \sum_{m=1}^{Nt} \sum_{n=1}^{Nt} \bar{\alpha}_m^i(s) D_{mn}^i \bar{\alpha}_n^i(s) \quad (\text{B-4})$$

$$D_{mn}^i = \frac{\pi(a^i)^2 (m-1)(n-1)}{2L(m+n-2)} \sum_{k=1}^{Ne} E_k^i * \left[ \left( e^{-(m+n-2)(z_k^i - \Delta t_k^i/2)} - e^{-(m+n-2)(z_k^i + \Delta t_k^i/2)} \right) \right] \quad (\text{B-5})$$

where  $z_k^i$  and  $\Delta t_k^i$  denote the distance from the top to the middle of the  $k^{th}$  element and the thickness of the  $k^{th}$  element of the  $i^{th}$  pile respectively.

## APPENDIX C

This appendix presents the relation between the matrices  $\beta^j$  and  $\omega^j$ . The relationship between unknown body forces and the vertical displacements of the extended half-space can be expressed as

$$\begin{bmatrix} \mathbf{F}^{1,1} & \mathbf{F}^{1,2} & \dots & \mathbf{F}^{1,j} & \dots & \mathbf{F}^{1,N_p} \\ \mathbf{F}^{2,1} & \mathbf{F}^{2,2} & \dots & \mathbf{F}^{2,j} & \dots & \mathbf{F}^{2,N_p} \\ \vdots & \vdots & \ddots & \vdots & \ddots & \vdots \\ \mathbf{F}^{i,1} & \mathbf{F}^{i,2} & \dots & \mathbf{F}^{i,j} & \dots & \mathbf{F}^{i,N_p} \\ \vdots & \vdots & \ddots & \vdots & \ddots & \vdots \\ \mathbf{F}^{N_p,1} & \mathbf{F}^{N_p,2} & \dots & \mathbf{F}^{N_p,j} & \dots & \mathbf{F}^{N_p,N_p} \end{bmatrix} \begin{Bmatrix} \mathbf{B}^1 \\ \mathbf{B}^2 \\ \vdots \\ \mathbf{B}^j \\ \vdots \\ \mathbf{B}^{N_p} \end{Bmatrix} = \begin{Bmatrix} \mathbf{w}^1 \\ \mathbf{w}^2 \\ \vdots \\ \mathbf{w}^j \\ \vdots \\ \mathbf{w}^{N_p} \end{Bmatrix} \quad (C-1)$$

or

$$\begin{aligned} \mathbf{F}^{1,1}\mathbf{B}^1 + \mathbf{F}^{1,2}\mathbf{B}^2 + \dots + \mathbf{F}^{1,j}\mathbf{B}^j + \dots + \mathbf{F}^{1,N_p}\mathbf{B}^{N_p} &= \mathbf{w}^1 \\ \mathbf{F}^{2,1}\mathbf{B}^1 + \mathbf{F}^{2,2}\mathbf{B}^2 + \dots + \mathbf{F}^{2,j}\mathbf{B}^j + \dots + \mathbf{F}^{2,N_p}\mathbf{B}^{N_p} &= \mathbf{w}^2 \\ &\vdots \\ \mathbf{F}^{i,1}\mathbf{B}^1 + \mathbf{F}^{i,2}\mathbf{B}^2 + \dots + \mathbf{F}^{i,j}\mathbf{B}^j + \dots + \mathbf{F}^{i,N_p}\mathbf{B}^{N_p} &= \mathbf{w}^j \\ &\vdots \\ \mathbf{F}^{N_p,1}\mathbf{B}^1 + \mathbf{F}^{N_p,2}\mathbf{B}^2 + \dots + \mathbf{F}^{N_p,j}\mathbf{B}^j + \dots + \mathbf{F}^{N_p,N_p}\mathbf{B}^{N_p} &= \mathbf{w}^{N_p} \end{aligned} \quad (C-2)$$

where  $\mathbf{F}^{i,j}$  denotes the Laplace transform of the vertical displacement of the  $i^{th}$  pile due to a vertical body force of the  $j^{th}$  pile and  $\mathbf{w}^j$  vertical displacement of the  $j^{th}$  pile. For each equation in the system in equation (C-2) yields

$$\sum_{i=1}^{N_p} \mathbf{F}^{i,1}\mathbf{B}^1 + \sum_{i=1}^{N_p} \mathbf{F}^{i,2}\mathbf{B}^2 + \dots + \sum_{i=1}^{N_p} \mathbf{F}^{i,j}\mathbf{B}^j + \dots + \sum_{i=1}^{N_p} \mathbf{F}^{i,N_p}\mathbf{B}^{N_p} = \sum_{j=1}^{N_p} \mathbf{w}^j \quad (C-3)$$

Substituting  $\mathbf{w}^j = \omega^j \alpha^j$  and  $\mathbf{B}^j = \beta^j \alpha^j$  into equation (C-3) results in

$$\sum_{i=1}^{N_p} \mathbf{F}^{i,1}\beta^1\alpha^1 + \sum_{i=1}^{N_p} \mathbf{F}^{i,2}\beta^2\alpha^2 + \dots + \sum_{i=1}^{N_p} \mathbf{F}^{i,j}\beta^j\alpha^j + \dots + \sum_{i=1}^{N_p} \mathbf{F}^{i,N_p}\beta^{N_p}\alpha^{N_p} = \sum_{j=1}^{N_p} \omega^j \alpha^j \quad (C-4)$$

The above equation can be rewritten in the following matrix form

$$\left[ \sum_{i=1}^{Np} \mathbf{F}^{i,1} \beta^1 \quad \sum_{i=1}^{Np} \mathbf{F}^{i,2} \beta^2 \quad \dots \quad \sum_{i=1}^{Np} \mathbf{F}^{i,j} \beta^j \quad \dots \quad \sum_{i=1}^{Np} \mathbf{F}^{i,Np} \beta^{Np} \right] \begin{Bmatrix} \alpha^1 \\ \alpha^2 \\ \vdots \\ \alpha^j \\ \vdots \\ \alpha^{Np} \end{Bmatrix} = \left[ \omega^1 \quad \omega^2 \quad \dots \quad \omega^j \quad \dots \quad \omega^{Np} \right] \begin{Bmatrix} \alpha^1 \\ \alpha^2 \\ \vdots \\ \alpha^j \\ \vdots \\ \alpha^{Np} \end{Bmatrix}$$

(C-5)

Finally, are obtained the following relations

$$\sum_{i=1}^{Np} \mathbf{F}^{i,j} \beta^j = \omega^j \quad (C-6)$$

## BIOGRAPHY

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