

## Chapter II

Phase Equilibrium Flash Vaporization  
and Partial CondensationTwo-Phase Flash Vaporization

A *flash* is a single-stage distillation in which a feed is partially vaporized to give a vapor that is richer in the more volatile components. In Fig.2.1a a liquid feed is

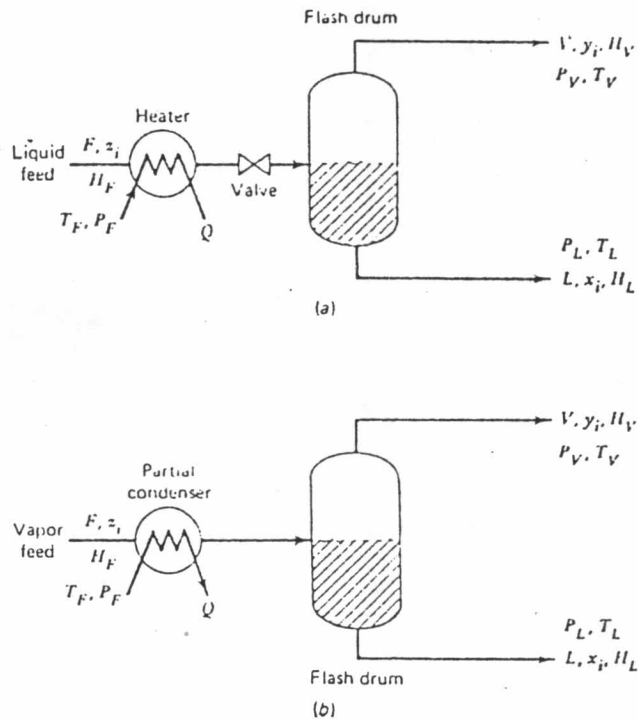


Figure 2.1 Continuous single-stage equilibrium separation. (a) Flash vaporization. (Adiabatic flash with valve; isothermal flash without valve when  $T_V$  is specified. (b) Partial condensation. (Analogous to isothermal flash when  $T_V$  is specified.

heated under pressure and flashed adiabatically across a valve to a lower pressure. The vapor being separated from the liquid residue in a flash drum if the valve is omitted, a low-pressure liquid can be partially vaporized in the heater and then separated into two phases. Alternatively, a vapor feed can be cooled and partially condensed, with phase separation in a flash drum as in Fig. 2.1b to give a liquid that is richer in the less volatile component. In both cases, if the component is properly designed, the vapor and liquid leaving the drum are in equilibrium. Unless the relative volatility is very large, the degree of separation achievable between two

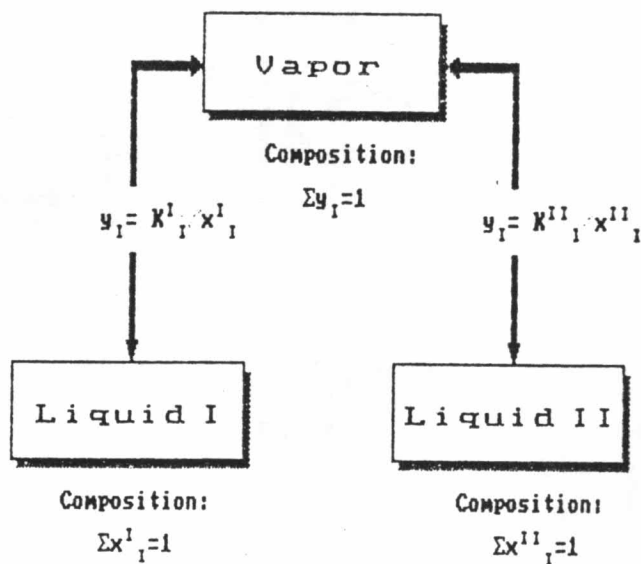


Figure 2.2

Equilibrium Relationships Between  
Immiscible Liquids and Vapor

components in a single stage is poor; so flashing and partial condensation are usually auxiliary operations used to prepare feed streams for further processing. However, the computational methods used in single-stage calculation are of fundamental importance.

### Three-Phase Flash Vapourization

Conventional two-phase flash calculation methods are often inadequate to handle the situation of a feed being flashed to form three phases: two liquid phases and a vapor phase. For every component in a three-phase system, there will exist two sets of liquid-vapor equilibrium data; one in association with each liquid phase. Compositions of the liquid phases may be calculated, since both liquid phases must be simultaneously in equilibrium with the single vapor phase.

### Vapour-Liquid-Liquid Equilibrium

For vapor-liquid-liquid equilibrium in terms of fugacities becomes

$$f_{1v} = f_{1L}^I = f_{1L}^{II} \quad (2.1)$$

To form an equilibrium ratio, fugacities are replaced by equivalent expressions involving mole fractions are:

$$f_{1v} = \delta_{1v} y_1 f_{1v} \quad (2.2)$$

$$f_{1L}^I = \delta_{1L}^I x_1^I f_{1L}^{\circ I} \quad (2.3)$$

$$f_{1L}^{II} = \delta_{1L}^{II} x_1^{II} f_{1L}^{\circ II} \quad (2.4)$$

or  $f_{1v} = \phi_{1v} y_1 P \quad (2.5)$

$$f_{1L}^I = \phi_{1L}^I x_{1L}^I P \quad (2.6)$$

$$f_{1L}^{II} = \phi_{1L}^{II} x_{1L}^{II} P \quad (2.7)$$

From equation (2.1) these equations can be rearranged for K-values in terms :

$$K_1^I = \phi_{1L}^I / \phi_{1v} = y_1 / x_1^I \quad (2.8)$$

$$K_1^{II} = \phi_{1L}^{II} / \phi_{1v} = y_1 / x_1^{II} \quad (2.9)$$

or  $K_1^I = (\delta_{1L}^I / \delta_{1v}) (f_{1L}^{\circ I} / f_{1v}^{\circ}) \quad (2.10)$

$$K_1^{II} = (\delta_{1L}^{II} / \delta_{1v}) (f_{1L}^{\circ II} / f_{1v}^{\circ}) \quad (2.11)$$

### Isothermal-Isobaric Flash

The material balance for a three phase flash is:

$$f_i = L^I x_i^I + L^{II} x_i^{II} + V y_i ; i=1,2,\dots,n. \quad (2.12)$$

The two liquid-phase mol fractions can be expressed in terms of the vapor-phase mol-fractions by means of the vapor-liquid equilibrium coefficients,  $K_1^I$  and  $K_1^{II}$ ;

$$f_i = L^I y_i / K_1^I + L^{II} y_i / K_1^{II} + V y_i \quad (2.13)$$

This equation can be solved for the vapor phase mol fractions;

$$y_1 = f_1 K_1' K_1'' / (VK_1' K_1'' + L^1 K_1'' + L'' K_1') \quad (2.14)$$

Alternately, all mol-fractions in equation(2.12) could be related to the mol-fraction in liquid-phase 1 or liquid-phase 2:

$$f_1 = L^1 x_1' + L'' K_1' x_1'' / K_1'' + VK_1' x_1' \quad (2.15)$$

$$f_1 = L^1 K_1'' x_1'' / K_1' + L'' x_1'' + VK_1'' x_1'' \quad (2.16)$$

equation then yield expressions for the mol-fractions in the liquid-phases:

$$x_1' = f_1 K_1'' / (VK_1' K_1'' + L^1 K_1'' + L'' K_1') \quad (2.17)$$

$$x_1'' = f_1 K_1' / (VK_1' K_1'' + L^1 K_1'' + L'' K_1') \quad (2.18)$$

Equation(2.14) may be summed over all components to yield:

$$1 = \sum_{i=1}^n (f_i K_1' K_1'' / (VK_1' K_1'' + L^1 K_1'' + L'' K_1')) \quad (2.19)$$

By an overall material balance:

$$V = F - L^1 - L'' \quad (2.20)$$

Inserting equation(2.20) into equation(2.19) and rearranging.

where:  $u_1 = L^1 / F$  and  $u_2 = L'' / F$

$$\begin{aligned}
P_1(u_1, u_2) &= \\
&= \sum_{i=1}^n K_1^i K_1^{i+1} / (K_1^i K_1^{i+1} + u_1 K_1^{i+1} (1 - K_1^i) + u_2 K_1^i (1 - K_1^{i+1})) - 1 \\
&= \sum_{i=1}^n (-u_1 K_1^{i+1} (1 - K_1^i) - u_2 K_1^i (1 - K_1^{i+1})) / (K_1^i K_1^{i+1} + \\
&\quad u_1 K_1^{i+1} (1 - K_1^i) + u_2 K_1^i (1 - K_1^{i+1})) \\
&= 0
\end{aligned} \tag{2.21}$$

Likewise, equations (2.17) and (2.18) become;

$$\begin{aligned}
P_2(u_1, u_2) &= \sum_{i=1}^n (1 - u_1) K_1^{i+1} (1 - K_1^i) - u_2 K_1^i (1 - K_1^{i+1}) / \\
&\quad (K_1^i K_1^{i+1} + u_1 K_1^{i+1} (1 - K_1^i) + u_2 K_1^i (1 - K_1^{i+1})) \\
&= 0
\end{aligned} \tag{2.22}$$

$$\begin{aligned}
P_3(u_1, u_2) &= \sum_{i=1}^n (-u_1 K_1^{i+1} (1 - K_1^i) + (1 - u_2) K_1^i (1 - K_1^{i+1})) / \\
&\quad (K_1^i K_1^{i+1} + u_1 K_1^{i+1} (1 - K_1^i) + u_2 K_1^i (1 - K_1^{i+1})) \\
&= 0
\end{aligned} \tag{2.23}$$

Two independent functions which have properties more suitable for numerical solution can be derived from equations (2.21-2.23). These are;

$$\begin{aligned}
 Q_1(u_1, u_2) &= P_2(u_1, u_2) - P_1(u_1, u_2) \\
 &= \sum_{i=1}^n z_i K_i^{11} (1 - K_i^{11}) / (K_i^{11} K_i^{11}) \\
 &\quad + u_1 K_i^{11} (1 - K_i^{11}) + u_2 K_i^{11} (1 - K_i^{11}) \quad (2.24)
 \end{aligned}$$

$$\begin{aligned}
 Q_2(u_1, u_2) &= P_3(u_1, u_2) - P_1(u_1, u_2) \\
 &= \sum_{i=1}^n z_i K_i^{11} (1 - K_i^{11}) / (K_i^{11} K_i^{11}) \\
 &\quad + u_1 K_i^{11} (1 - K_i^{11}) + u_2 K_i^{11} (1 - K_i^{11}) \quad (2.25)
 \end{aligned}$$

The  $P$ -functions can be written in terms of the  $Q$ -functions:

$$P_1(u_1, u_2) = -u_1 Q_1(u_1, u_2) - u_2 Q_2(u_1, u_2) \quad (2.26)$$

$$P_2(u_1, u_2) = (1-u_1) Q_1(u_1, u_2) - u_2 Q_2(u_1, u_2) \quad (2.27)$$

$$P_3(u_1, u_2) = -u_1 Q_1(u_1, u_2) - (1-u_2) Q_2(u_1, u_2) \quad (2.28)$$

Simultaneous solution of equation(2.24,2.25)will yield the same values of  $u_1$  and  $u_2$  as simultaneous solution of equation(2.21-2.23). But the  $Q$ -function are more suitable for iterative solution because they contain no superfluous roots in the domain of interest. Once it has been determined that three phases are present, a (2\*2 matrix) search must be performed for the values of  $u_1$  and  $u_2$  which cause equations(2.24,2.25) to be satisfied. If the flash is being performed with recalculated  $K$ -values from a previously

converged solution, the previous value of  $u_1$  and  $u_2$  are good starting estimates.

The Newton-Raphson iteration procedure can be used to search for the solution. The required partial derivatives of  $Q_1$  and  $Q_2$  with respect to  $u_1$  and  $u_2$  are

$$\begin{aligned} \frac{dQ_1}{du_1} = & \sum_{i=1}^n (-z_i K_i^{(1)} K_i^{(1)} (1-K_i^{(1)})^2) / (K_i^{(1)} K_i^{(1)} + u_1 K_i^{(1)} (1-K_i^{(1)})) \\ & + u_2 K_i^{(1)} (1-K_i^{(1)})^2 \end{aligned} \quad (2.29)$$

$$\begin{aligned} \frac{dQ_1}{du_2} = \frac{dQ_2}{du_1} = & \sum_{i=1}^n (-z_i K_i^{(1)} K_i^{(1)} (1-K_i^{(1)}) (1-K_i^{(1)})) / (K_i^{(1)} K_i^{(1)} \\ & + u_1 K_i^{(1)} (1-K_i^{(1)}) + u_2 K_i^{(1)} (1-K_i^{(1)})^2) \end{aligned} \quad (2.30)$$

$$\begin{aligned} \frac{dQ_2}{du_2} = & \sum_{i=1}^n (-z_i K_i^{(1)} K_i^{(1)} (1-K_i^{(1)})^2) / (K_i^{(1)} K_i^{(1)} + u_1 K_i^{(1)} (1-K_i^{(1)})) \\ & + u_2 K_i^{(1)} (1-K_i^{(1)})^2 \end{aligned} \quad (2.31)$$

The solution procedure consists of finding successive estimates of the root by the iterative application of

$$\begin{array}{c|c} u_1 \\ \hline \end{array} \Big|_{m+1} = \begin{array}{c|c} u_1 \\ \hline \end{array} \Big|_m + \begin{array}{c|c|c|c} \frac{dQ_1}{du_1} & \frac{dQ_1}{du_2} & & \\ \hline & & & \\ \frac{dQ_2}{du_1} & \frac{dQ_2}{du_2} & & \\ \hline & & & \\ Q_1 & & & \\ \hline & & & \\ Q_2 & & & \\ \hline & & & \end{array} \Big|_m$$



$$= \begin{vmatrix} u_1 & & dQ_2/du_2 & -dQ_1/du_2 \\ & - (1/D) & & \\ u_2 & & -dQ_2/du_1 & dQ_1/du_1 \end{vmatrix} \begin{vmatrix} Q_1 \\ \\ Q_2 \end{vmatrix} \quad (2.32)$$

where  $D$  is the determinant of the Jacobian matrix

$$D = dQ_1/du_1 \cdot dQ_2/du_2 - dQ_1/du_2 \cdot dQ_2/du_1 \quad (2.33)$$

Convergence is declared when  $Q_1$  and  $Q_2$  are zero to within tolerance.

### Criteria of Phase Equilibria

For each phase in a multiphase, multicomponent system, the Gibbs free energy is given functionally as

$$G = G(T, P, n_1, n_2, \dots, n_c) \quad (2.34)$$

where  $n$  = moles and subscripts refer to species. The total differential of  $G$  is

$$dG = \left( \frac{dG}{dT} \right)_{P, n} dT + \left( \frac{dG}{dP} \right)_{T, n} dP + \sum_{i=1}^c \left( \frac{dG}{dn_i} \right)_{P, T, n} dn_i \quad (2.35)$$

where  $j = 1$ . From classical thermodynamics

$$\left( \frac{dG}{dT} \right)_{P, n} = -S \quad (2.36)$$

and 
$$\left(\frac{dG}{dP}\right)_{T,n} = V \quad (2.37)$$

where  $S =$  entropy and  $V =$  volume. Defining the *chemical potential*,  $\mu$ , of species  $i$  as

$$\mu_i = \left(\frac{dG}{dn_i}\right)_{P,T,n} \quad (2.38)$$

and substituting into (2.35), we have

$$dG = -S dT + V dP + \sum_{i=1}^c \mu_i dn_i \quad (2.39)$$

When (2.39) is applied to a closed system consisting of two phases in equilibrium at uniform temperature and pressure, where each phase is an open system capable of mass transfer with another phase

$$dG_{\text{system}} = \sum_{k=1}^p \left( \sum_{i=1}^c \mu_i^k dn_i^k \right) = 0 \quad (2.40)$$

where the superscript (k) refers to each of  $p$  phases. Conservation of moles of each species requires that

$$dn_1^1 = - \sum_{k=2}^p dn_1^k$$

which, upon substitution into (2.39), gives

$$\sum_{k=2}^p \left( \sum_{i=1}^c (\mu_i^k - \mu_i^1) dn_i^k \right) = 0 \quad (2.41)$$

with  $dn_1^1$  eliminated in (2.41), each  $dn_1^k$  term can be varied independently of any other  $dn_1^k$  term. But this requires that each coefficient of  $dn_1^k$  in (2.41) be zero. Therefore,

$$\mu_1^1 = \mu_1^2 = \mu_1^3 = \dots = \mu_1^p \quad (2.42)$$

Thus, the chemical potentials of any species in a multicomponent system are identical in all phases at physical equilibrium. Chemical potential cannot be expressed as an absolute quantity, and the numerical values of chemical potential are difficult to more easily understood physical quantities. Furthermore, the chemical potential approaches an infinite negative value as pressure approaches zero. For these reasons, the chemical potential is not directly useful for phase equilibria calculations. Instead, fugacity, as defined below, is employed, as a surrogate. Equation (2.38) restated in terms of chemical potential is

$$(d\mu_1/dP)_T = \bar{v}_1 \quad (2.43)$$

where  $\bar{v}_1$  = partial molal volume. For a pure substance that behaves as an ideal gas

$$\bar{v}_1 = RT/P$$

and (2.43) can be integrated to give

$$\mu_1 = RT \ln P + C_1(T) \quad (2.44)$$

where  $C_1$  depend on  $T$ . Unfortunately, (2.44) does not describe real multicomponent gas or liquid behavior. However, (2.44) was rescued by G.M. Lewis, who in 1901 invented the fugacity  $f$ , a pseudopressure when used in place of pressure in (2.44), thus for a component in a mixture.

$$f_1 = C_2(T) \exp(\mu_1/RT) \quad (2.45)$$

where  $C_2$  is related to  $C_1$ . Regardless of the value of  $C_1$  at physical equilibrium, (2.42) can be replaced with

$$f_1^1 = f_1^2 = f_1^3 = \dots f_1^p \quad (2.46)$$

For a pure, ideal gas, fugacity is equal to the pressure and, for a component in an ideal gas mixture, it is equal to its partial pressure,  $p_1 = y_1 P$ .