

แถบอสุทธีบริ เวศทางในสารกึ่งตัวนำที่ถูกโคป.

อย่างหนัก โดยวิธีนับต่ำสุด



นาย เผด็จชัย ไชยสิทธิ์

004186

วิทยานิพนธ์นี้ เป็นส่วนหนึ่งของการศึกษาตามหลักสูตรปริญญาวิทยาศาสตรมหาบัณฑิต

ภาควิชา เคมี

บัณฑิตวิทยาลัย จุฬาลงกรณ์มหาวิทยาลัย

พ.ศ. ๒๕๒๕

I 16580618

IMPURITY-BAND TAILS IN HEAVILY DOPED
SEMICONDUCTORS : MINIMUM COUNTING METHODS

Mr. Pachernchai Chaiyasith

A Thesis Submitted in Partial Fulfillment of the Requirements

for the Degree of Master of Science

Department of Chemistry

Graduate School

Chulalongkorn University

1982

เมื่อ $y = \beta/Q$ และ v ซึ่งเขียนอยู่ในหน่วยของ $E_Q = \hbar^2 Q^2/2m$ คือพลังงานส่วนที่อยู่ต่ำกว่าศักย์มีขั้ว E_0 นอกจากนี้ยังคำนวณและเปรียบเทียบผลเชิงตัวเลขที่ได้กับผลของฮิลเปอรินกับลักษณะและผลของอิมารคกับตุราฟออร์ก ซึ่งพบว่า ผลที่ได้สอดคล้องกับผลการคำนวณของฮิลเปอรินกับลักษณะดีกว่าผลของอิมารคกับตุราฟออร์ก

สำหรับศักย์อสุธิแบบ เกาส์ ฟังก์ชันเชิงทดลองที่ใช้ สมมติว่าเป็นฟังก์ชันคลื่นสถานะพื้นแบบ เกาส์ $f(\vec{x}) = (2\xi_0/\pi)^{3/4} \exp(-\xi_0^2 x^2)$ เมื่อ $\xi_0 = m\omega/2\hbar$ ซึ่งทำให้ความหนาแน่นสถานะส่วนทาง เขียนอยู่ในนิพจน์เชิงวิเคราะห์ได้ และเหมือนกันทุกประการ กับที่ได้จากทฤษฎีของศาสตราจารย์ ดร.วิรุทธิ์ สายคณิต ที่ใช้วิธีอินทิเกรตตามทางของพายน์แมน นั่นคือ

$$\rho(E) = [(E_L/L)/\xi_L^2] a(v) \exp[-E_L^2 b(v)/2\xi_L]$$

$$\text{เมื่อ } a(v) = [(1 + 16v)^{1/2} - 1]^{3/2} [(1 + 16v)^{1/2} + 7]^{9/2} / 2^{12} 2^{1/2} \pi^2$$

$$\text{และ } b(v) = [(1 + 16v)^{1/2} - 1]^{1/2} [(1 + 16v)^{1/2} + 7]^{7/2} / 2^8$$

เมื่อ v ซึ่งเขียนอยู่ในหน่วยของ $E_L = \hbar^2/2mL^2$ คือพลังงานส่วนที่อยู่ต่ำกว่าศักย์มีขั้ว E_0

Thesis Title Impurity-Band Tails in Heavily Doped Semiconductors :
 Minimum Counting Methods

Name Mr. Pachernchai Chaiyasith

Thesis Advisor Assistant Professor Sirirat Kokpol, Ph.D.
 Professor Virulh Sa-yakanit, F.D.

Department Chemistry

Academic Year 1981

ABSTRACT

A simplified approach of minimum counting method (Halperin and Lax Theory) of the density of states tail in heavily doped semiconductors is presented.

For a screened Coulomb potential the trial wave function is assumed to be in the form $\phi(r) = (2\beta^4/3\pi)^{1/2} r^{1/2} \exp(-\beta r)$. The density of states tail can be expressed in analytical form

$$\rho(E) = [(E_Q/Q)^3 / \epsilon^2] a(v) \exp [-E_Q^2 b(v) / 2\epsilon]$$

$$\text{where } a(v) = \frac{546 (2v+y^2)^3 (2y+1)^{14} (264y^3 + 60y^2 + 14y+1)^{3/2}}{(27)^{3/2} \pi^2 y^3 (1152y^4 + 728y^3 + 244y^2 + 42y+3)^{7/2}}$$

$$\text{and } b(v) = \frac{9 (2v+y^2)^2 (2y+1)^7}{4y^3 (1152y^4 + 728y^3 + 244y^2 + 42y+3)}$$

with $y = \beta/Q$ and v being the energy below the mean potential E_0 in units of $E_Q = \hbar^2 Q^2/2m$. Numerical results are calculated and compared with Halperin and Lax and with Eymard and Duraffourg results. It is found that, the results are in better agreement with computed Halperin and Lax results than Eymard and Duraffourg results.

For a Gaussian impurity potential, the trial wave function is assumed to be the ground state Gaussian wave function,
 $f(\vec{x}) = (2\xi_0/\pi)^{3/4} \exp(-\xi_0^2 x^2)$ where $\xi_0 = m\omega/2\hbar$. An analytical expression of the density of states tail is obtained, identically with Sa-yakanit theory using Feynman path integral method, i.e.

$$\rho(E) = [(E_L/L)^3 / \xi_L^2] a(v) \exp[-E_L^2 b(v) / 2\xi_L]$$

$$\text{where } a(v) = [(1 + 16v)^{1/2} - 1]^{3/2} [(1 + 16v)^{1/2} + 7]^{9/2} / 2^{12} \pi^{1/2} 2^2$$

$$\text{and } b(v) = [(1 + 16v)^{1/2} - 1]^{1/2} [(1 + 16v)^{1/2} + 7]^{7/2} / 2^8, \text{ with}$$

v being the energy below the mean potential E_0 in units of $E_L = \hbar^2/2mL^2$.

ACKNOWLEDGEMENTS



The author wishes to express his deep gratitude to Assistant Professor Dr. Sirirat Kokpol and Professor Dr. Virulh Sa-yakanit for their helpful guiding, advising and encouraging throughout the course of this research. He is grateful to Archan Chai Hok Eab and Archan Noppadon Suttisiri for their valuable suggestions. Grateful acknowledgement is also accorded to Archan Rangsan Chalurmsri for his advising in writing computer program. Finally, the author also wishes to thank the University Development Commission for granting a scholarship.

CONTENTS

	PAGE
ABSTRACT (in Thai)	iv
ABSTRACT	vi
ACKNOWLEDGEMENTS	viii
LIST OF TABLES	xi
LIST OF FIGURES	xiv
CHAPTER	
I. INTRODUCTION	
1.1 Perfect Crystals	2
1.2 Imperfect Crystals	7
1.3 Heavily Doped Semiconductors	9
II QUANTUM THEORY OF IMPURITY-BAND TAILS	
2.1 Halperin and Lax Theory (Minimum Counting Method)	15
2.1.1 Minimum Counting Method	15
2.1.2 Density of States	19
2.1.3 Screened Coulomb Potential	28
2.2 Sa-yakanit Theory	33
2.2.1 Density of States	33
2.2.2 Gaussian Potential	37



	PAGE
2.2.3 Screened Coulomb Potential	42
III THE SIMPLIFIED APPROACH OF THE HALPERIN AND LAX THEORY	
3.1 Screened Coulomb Potential.....	48
3.1.1 Eymard and Duraffourg Method.....	49
3.1.2 Present Method.....	61
3.2 Gaussian Potential.....	72
3.2.1 Present Method.....	72
IV CONCLUSION AND DISCUSSION	
4.1 Conclusion.....	84
4.2 Comparison of Results.....	87
4.2.1 Screened Coulomb Potential.....	87
4.2.2 Gaussian Potential.....	98
4.2.3 The Application of The Impurity Potential Models.....	103
REFERENCES	104
APPENDIX A FORTRAN IV PROGRAM FOR CALCULATING THE DENSITY OF STATES IN EYMARD AND DURAFFOURG METHOD.....	110
APPENDIX B FORTRAN IV PROGRAM FOR CALCULATING THE DENSITY OF STATES IN PRESENT METHOD.....	116
APPENDIX C FORTRAN IV PROGRAM FOR CALCULATING THE DENSITY OF STATES FOR THE CASE OF GAUSSIAN IMPURITY POTENTIAL IN PRESENT METHOD.....	120
VITA.....	122

LIST OF TABLES

TABLE		PAGE
2.1	The limiting values of $a(v)$, $b(v)$, and $n(v)$ for a screened Coulomb potential obtained from Halperin and Lax Theory.....	32
3.1	Numerical results of the functions $T(v)$, $n(v)$ and of the adjustable parameter x obtained from Eymard and Duraffourg method.....	57
3.2	Numerical results of the dimensionless functions, $a(v)$ and $b(v)$ obtained from Eymard and Duraffourg method.....	58
3.3	Numerical results of the functions, σ_0^2 , σ_1^2 , and μ obtained from Eymard and Duraffourg method.....	59
3.4	Numerical results of the density of states for $\xi' = 0.05, 0.5, 5, \text{ and } 50$ obtained from Eymard and Duraffourg method.....	60
3.5	The limiting values of $a(v)$, $b(v)$, $n(v)$, and $T(v)/v$ obtained from Eymard and Duraffourg method...67	67
3.6	Numerical results of the functions $T(v)$, $n(v)$, and the adjustable parameter y obtained from present method.....	68

LIST OF TABLES (continue)

TABLE		PAGE
3.7	Numerical results of the dimensionless functions $a(v)$ and $b(v)$ obtained from present method.....	69
3.8	Numerical results of the functions σ_0^2 , σ_1^2 and μ obtained from present method.....	70
3.9	Numerical results of the density of states for $\xi' = 0.05, 0.5, 5, \text{ and } 50$ obtained from present method.....	71
3.10	The limiting values of the function $a(v)$, $b(v)$, $n(v)$, and $T(v)/v$ obtained from present method..	81
3.11	Numerical results of the functions $T(v)$, $n(v)$ and of the adjustable parameter z (for Gaussian potential) obtained from present method.....	82
3.12	Numerical results of the dimensionless functions, $a(v)$ and $b(v)$ (for Gaussian potential) obtained from Present method.....	83
4.1	Comparison between the values of $a(v)$ obtained from four methods for a screened Coulomb potential....	89

LIST OF TABLES (continue)

TABLE		PAGE
4.2	Comparison between the values of $b(v)$ obtained from four methods for a screened Coulomb potential.....	91
4.3	Comparison between the limiting values of $a(v)$, $b(v)$, and $T(v)/v$ obtained from four methods for a screened Coulomb potential.....	93

LIST OF FIGURES

FIGURE		PAGE
1.1	The one-dimensional perfect crystal potential $V(x)$ as a function of position x	3
1.2	Energy band scheme of perfect crystal.....	5
1.3	The density of states $\rho(E)$ as a function of energy E	10
2.1	Potential wells in heavily doped semiconductor.....	19
3.1	The wave function $S(r)$ as a function of r in dimensionless unit $S(r) = 0.1$	62
4.1	The density of states for $2\xi' = 0.1$	94
4.2	The density of states for $2\xi' = 1$	94
4.3	The density of states for $2\xi' = 10$	95
4.4	The density of states for $2\xi' = 100$	95
4.5	The logarithmic derivative $n(\nu)$ of the exponent $b(\nu)$ as a function of energy ν	96
4.6	The density of states for $2\xi'_L = 0.1$	100

LIST OF FIGURES (continue)

FIGURE		PAGE
4.7	The density of states for $2\xi'_L = 1$	100
4.8	The density of states for $2\xi'_L = 10$	101
4.9	The density of states for $2\xi'_L = 100$	101
4.10	The logarithmic derivative $n(\nu)$ of the exponent $b(\nu)$ as a function of ν	102